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NOTION OF A JET IN A CROSSFLOW

Richard L. Fearn and Robert P. Weston

Langley Research Center

Hampton, Va. 23665



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16. Abstract				
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INDUCED PRESSURE DISTRIBUTION OF A JET IN A CROSSFLOW

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SUMMARY

An experimental investigation of the turbulent flow of a subsonic round jet exhausting perpendicularly from a flat plate into a subsonic crosswind of the same temperature has been conducted in the Langley V/STOL tunnel. The large test section of this tunnel made it possible to use a jet with the relatively large diameter of 10.16 cm (4.00 in.). Pressures were measured on the flat plate at over 400 locations to provide a detailed description of the static pressure distribution. Results are presented in tabular and graphical forms for jet to crossflow velocity ratios ranging from 2 to 10, and comparisons are made with results of other experiments. The results indicate that the dominant flow parameter affecting the pressure distributions is the ratio of jet to crossflow velocities. The distribution of turbulence on the plate also appears to be a function of velocity ratio.

INTRODUCTION

In the transition between hover and wingborne flight, a significant part of the lift of a VTOL aircraft is furnished by direct engine thrust. This injection of a high velocity jet of air at almost right angles to the crossflow caused by the aircraft's forward motion produces a complicated flow field which affects the aerodynamic characteristics of the aircraft. Several wind-tunnel tests have been conducted to measure these jet-induced effects on the pressure loading and aerodynamic coefficients for specific aircraft configurations (refs. 1 to 9).

In an attempt to describe in general the essential features of this complicated flow field, many investigators have concentrated on a simplified version of the situation and studied the turbulent flow of a subsonic round jet exhausting perpendicularly through a large flat plate into a subsonic crossflow. Such a study has application to areas other than V/STOL aerodynamics such as the process of cooling combustion gases in gas-turbine combustors or the discharge of an effluent into a waterway. For several decades experiments have been conducted on certain aspects of this jet in a crossflow problem, namely

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the path of the jets (refs. 10 to 16) and the pressure distribution on the flat plate (refs. 15 and 17 to 23). However, the pair of contrarotating vortices, which constitutes one of the dominant features of the velocity field, has only recently received detailed attention (refs. 15, 16, 24, and 25). Due to limitations in wind-tunnel size, most previous experiments have been restricted to jets with rather small diameters of 0.6 to 2.5 cm (0.25 to 1.0 in.) to avoid wind-tunnel wall effects. An extensive bibliography for work done prior to 1969 is given by R. J. Margason in reference 26.

An experimental investigation of the pressure and velocity fields associated with a jet in a crossflow for flows of the same temperature was conducted in the Langley V/STOL tunnel whose large test section made it feasible to use a jet with the relatively large diameter of 10.16 cm (4.0 in.). The investigation included measurements of the pressure distribution on the flat plate and the velocity field in the jet and resulted in a quantitative description of the pair of contrarotating vortices associated with the jet (ref. 25). The measured pressure distributions on the flat plate are presented in this report in tabular and graphical form for jet to crossflow velocity ratios ranging from 2 to 10.

A comparison of the results of the present experiment with those of a similar experiment conducted at the Office National d'Études et de Recherches Aérospatiales (ONERA) (refs. 21 and 22) is given also. This comparison is significant because the ONERA experiment is the only other one to date which utilized a jet of comparable size and is included herein because the results have heretofore received only limited distribution. Comparison is also made with the results of experiments using smaller jets (refs. 15, 17, 18, 20, and 23).

SYMBOLS

Values are given both in SI Units and U.S. Customary Units. The measurements and calculations were made in U.S. Customary Units.

- A area of jet orifice, m^2 (ft²)
- A_j jet exit area corrected for nozzle boundary-layer thickness, m^2 (ft²)
- C_p pressure coefficient, $\frac{p p_{\infty}}{q_{\infty}}$
- \overline{C}_p time average of C_p at a point
- C_p' fluctuations of C_p about its time-average value, \overline{C}_p
- D jet diameter, cm (in.)

Н height of plate above tunnel floor in jet diameters M Mach number m mass flow rate in jet nozzle, kg/sec (slugs/sec) Reynolds number based on jet diameter and crossflow velocity N_{Re} integer n static pressure, N/m^2 (lb/ft²) p total pressure in jet plenum, N/m² (lb/ft²) p_t dynamic pressure, N/m² (lb/ft²) \mathbf{q} effective velocity ratio defined by equation (1) \mathbf{R} gas constant, m-N/kg-K (ft-lb/slug-OR) Ē radial distance from origin of X,Y coordinate system, cm (in.) r temperature, K (OR) т U velocity, m/sec (ft/sec) coordinate system in plane of flat plate (see fig. 1(b)) X,Yratio of specific heat at constant pressure to specific heat at constant γ volume for a gas angle measured from negative X-axis, degrees (see fig. 1) θ kinematic viscosity, m²/sec (ft²/sec) mass density, kg/m³ (slugs/ft³)

standard deviation of fluctuations in Cn

σ

Subscripts:

- j denotes static property of jet fluid at nozzle-exit plane
- t denotes stagnation property of fluid in jet plenum
- ∞ denotes static property of crossflow fluid

APPARATUS

This experiment was conducted in the Langley V/STOL wind tunnel which is a closed return atmospheric tunnel with a test section 4.42 by 6.63 m (14.50 by 21.75 ft). Test-section airspeeds for this investigation ranged from 23 to 66 m/sec (75 to 215 ft/sec), or free-stream Mach numbers of 0.07 to 0.19. Jet Mach numbers ranged from 0.09 to 0.95.

The flat plate was originally constructed from a 1.22-m or 12D by 2.44-m or 24D (4.0- by 8.0-ft) sheet of 1.91-cm (0.75-in.) plywood with an overlay of 1.59-mm-thick (0.0625-in.) formica on the upper surface. A hole for the nozzle was centered midway between the long sides of the plate and 6D aft of its leading edge. The plate was mounted 3D above the wind-tunnel floor to remove the plate from the tunnel boundary layer, which was about 1.5D thick. In addition, boundary-layer suction could be applied to the tunnel floor upstream of the plate to reduce the tunnel-boundary-layer thickness.

Four hundred and fifty-nine static pressure orifices were placed on the original flat plate at selected intersections of ray and circles centered at the jet orifice as indicated in figure 1(a). Each static pressure port was made of 1.02-mm (0.040-in.) inside diameter stainless-steel thin-wall tubing mounted flush with the upper surface of the flat plate and protruding several centimeters through the lower surface. Sheet metal fairing was used on the underside of the plate to protect the static pressure leads and to streamline the nozzle-plenum assembly.

The majority of the data presented in this report was taken using a 24D by 27D plate resulting from front and side extensions to the original plate (fig. 1(b)) and with the plate mounted 4.5D above the wind-tunnel floor. For both sizes of the flat plate, the leading edge was rounded with a radius of about 1.02 cm (0.40 in.) and a transition strip was placed 3.8 cm (1.5 in.) aft of the leading edge.

The jet of air was formed by using a plenum chamber and a 20:1 convergent nozzle (ref. 27) designed to provide a flat velocity profile at the 10.16-cm-diameter (4.00-in.) nozzle exit. In order to measure the static pressure in the jet as close to the plane of the plate as possible, a small ring (0.125D) was installed at the jet orifice with four static pressure ports spaced equally around the inside surface of this jet exit ring 0.635 cm

(0.25 in.) from the exit plane. When the plate was mounted 4.5D above the tunnel floor, a constant diameter cylindrical extension was used to extend from the convergent nozzle to the jet exit ring. Figure 2 is a sketch of the nozzle-plenum assembly with the extension in place, and figure 3 is a photograph of the plate and jet orifice in the wind tunnel. Supply air for the jet was heated so that the temperature at the jet orifice would be approximately the same as that of the crossflow. A Venturi flowmeter in the air supply line was used to measure the mass flow rate of the jet.

All pressures were measured with pressure transducers which were calibrated with water or mercury manometers. The leads from the static pressure ports on the flat plate and from the boundary-layer mouse, which was used to measure the boundary layer on the flat plate, were connected by plastic tubing to 11 pressure scanning devices mounted on the underside of the plate. Each device consisted of a single pressure transducer which could be connected sequentially to each of 48 pressure tubes. The output of each pressure transducer was fed into a r.c. low pass filter to attenuate fluctuations in the transducer output signal. One second after each static pressure port was connected to the pressure transducer, the signal from the filter was sampled and recorded on magnetic tape.

TEST PROCEDURES AND CONDITIONS

Jet-Nozzle Characteristics

Prior to the wind-tunnel test, the jet-nozzle characteristics were investigated with no crossflow. A total pressure probe and a static pressure probe were used for surveys in the nozzle-exit plane and in several parallel planes up to 20D from the exit plane. These pressures were measured with water and with mercury manometers. The turbulent boundary-layer displacement thickness of the jet at the nozzle-exit plane (without extensions) was determined to be $(2.8 \pm 0.5) \times 10^{-3}$ D. Outside of the nozzle boundary layer the exit-plane velocity profile was flat to within 0.5 percent of the velocity measured at the center of the jet. The jet core extended almost 6D from the exit plane. The jet fluid at the exit had low turbulence, but once outside the jet core, fluid mixing created large turbulence in the jet.

Test Conditions

The investigation of the pressure distribution on the flat plate was conducted during three separate wind-tunnel tests. In the first test, the original flat plate (12D by 24D) was mounted 3D above the tunnel floor. Results of this experiment caused doubts about the accuracy of the measured pressures near the leading and side edges of the plate. It was decided to repeat the test to study the effect on the pressure distribution of changing various test conditions. The diagnostic phase of the second test included the test conditions shown in table 1.

TABLE 1.- TEST CONDITIONS

Н	Extensions to flat plate	Boundary-layer suction on wind-tunnel floor
3.0	None	On
3.0	None	Off
3.0	Leading edge	On
3.0	Leading edge	Off
4.5	Leading edge and side	On
4.5	Leading edge and side	Off

The remainder of the second test was made with the plate 4.5D above the floor, all extensions in place, and with no boundary-layer suction on the wind-tunnel floor. The pressure distribution was measured for a range of effective jet to crossflow velocity ratios, holding the jet Mach number constant. This was repeated for jet Mach numbers of 0.50, 0.75, and 0.93. Data for comparison with ONERA (ref. 21) were taken by matching wind-tunnel and jet Mach numbers as well as the effective velocity ratios. For comparison with the results of other experiments (refs. 15, 17, 18, and 23), only the effective velocity ratios were matched.

During a third wind-tunnel test, the pressure at selected ports was measured as a function of time for effective velocity ratios of 4 and 8 in an attempt to understand the nature of scatter in the data of the previous tests for certain locations on the plate. This was accomplished by sampling the output of the transducer connected to a given static pressure port for an interval of 10 seconds at a sampling rate of 31 times per second. The data were taken both with and without the low pass filters in the circuit. In this test, the plate was mounted 3D above the tunnel floor and consisted of the original plate with the leading-edge extension only. This configuration was used since the results of the diagnostic tests showed the leading-edge extension to be the only significant change in the conditions.

Effective Velocity Ratio

The properties of a jet in a crossflow depend primarily on the ratio of the momentum flux across the jet orifice to the momentum flux of the crossflow over an equal area. It is useful to define an effective velocity ratio as the square root of this ratio of momentum fluxes

$$R = \left(\frac{\int_{A} \rho_{j} U_{j}^{2} dA}{\rho_{\infty} U_{\infty}^{2} A}\right)^{1/2}$$
(1)

The reciprocal of equation (1) is used by some investigators and it is usually referred to by the same name.

If the jet and crossflow fluids are the same and obey the ideal gas equation of state, and if the speed U_j is constant across the jet orifice, then the effective velocity ratio may be written: $R = \frac{M_j}{M_\infty} \left(\frac{p_j}{p_\infty}\right)^{1/2}.$ In the experiment, this equation could be used in four ways to calculate the effective velocity ratio. The first involves measurements of the mass flow rate of the jet, the total temperature and total pressure of the jet, and the area of the jet orifice. Using the nozzle boundary-layer displacement thickness measured with no crossflow to give an estimate of the jet exit area A_j , the jet Mach number can be calculated, assuming isentropic flow, from $M_j \left(1 + 0.2 M_j^2\right)^{-3} = \frac{\dot{m}}{p_t A_j} \left(\frac{R}{r} T_t\right)^{1/2}$. With M_j known, the jet-exit-plane pressure can be calculated from $p_j = p_t \left(1 + 0.2 M_j^2\right)^{-3.5}$. A second way of using the equation for the effective velocity ratio makes use of the average pressure measured by the jet exit ring to give the static pressure at the jet exit plane. Isentropic expansion to this average pressure gives the jet Mach number from the relation $M_j = \left\{ \int \left[\frac{p_t}{p_j}\right]^{2/7} - 1 \right\}^{1/2}$. In reality, the pressures at the four jet-exit-ring ports differed greatly and it is questionable whether their average pressure actually represents the jet pressure at the exit.

If the jet static pressure at the exit plane is assumed equal to the crossflow static pressure, the equation for the effective velocity ratio simplifies to a ratio of the jet to crossflow Mach numbers. Additionally, if the jet and crossflow temperatures are equal, this becomes the actual velocity ratio, $R = \frac{Uj}{U_{\infty}}$. For use in the equation $R = \frac{Mj}{M_{\infty}}$, the jet Mach number can be calculated by two methods: using (1) the mass flow rate of the jet or (2) the jet total pressure. The isentropic relations for these are

$$M_j(\dot{m}) = \frac{\dot{m}}{p_{\infty}A_i} \left(\frac{\overline{R}T_j}{\gamma}\right)^{1/2}$$

or

$$\mathbf{M}_{j}(\mathbf{p}_{t}) = \left\{ 5 \left[\left(\frac{\mathbf{p}_{t}}{\mathbf{p}_{\infty}} \right)^{2/7} - 1 \right] \right\}^{1/2}$$

Differences in jet Mach number of up to 0.08 and in R of up to 14 percent are obtained between calculations using expansion to average jet-exit-ring pressure and expansion to crossflow static pressure. For presentation of results in this report the calculations are based on expansion to crossflow static pressure because of uncertainties in the

jet exit pressures due to difficulties in their measurement. In addition, calculating R from expansion to crossflow static pressure provides consistent comparison with data from the other investigators who calculate their velocity ratios in a similar way.

Flat-Plate Boundary Layer

The boundary layer on the flat plate was measured at several locations along the upstream ray from the jet for each test condition (see table 1). It was found that changing the test conditions had little effect on the behavior of the boundary layer on the plate. There was a separation bubble at the rounded leading edge of the plate and reattachment of the flow as a thickened turbulent boundary layer. Whether the jet was on or off did not significantly affect the measured boundary-layer velocity profiles. The estimated displacement thickness of the boundary layer at the location of the jet was 0.07D, based on boundary-layer measurements upstream of the jet.

PRESENTATION OF RESULTS

The pressure distribution on the flat plate is presented in terms of a pressure coefficient $C_p = \frac{p-p_\infty}{q_\infty}$, where p is the measured static pressure at a specific location on the plate. Except when otherwise noted, p is the difference between the pressure measured at the test condition $\left(R = \frac{M_j}{M_\infty}\right)$ and the pressure measured with the jet off $\left(M_j = 0\right)$, and with the tunnel Mach number unchanged.

The results of the wind-tunnel investigation of the pressure distribution of the flat plate are presented as follows:

	Figure	Table
Values of C_p for a range of R	4 to 11	
Time variation for C_p at selected locations for $R=4.1$ and $R=8$.		2
Uncertainty bands for contour plots	17	
Effect of changing jet Mach number holding R constant	18	
Comparison with the results of other experiments	19 to 23	3

RESULTS AND DISCUSSION

Pressure Distribution on the Flat Plate

Detailed descriptions of the measured pressure distribution on the flat plate are presented in figures 4 to 11 for effective velocity ratios of 2.2, 2.8, 3.9, 5.1, 6.1, 7.0, 8.0, and 10.0, respectively. Each figure presents the pressure distribution for one effective velocity ratio in four different formats. Numerical values of C_p are presented in tabular form in part (a). The pressure distribution as a function of radial distance from the jet for each ray of pressure ports is shown in part (b), and in part (c) it is shown as a

function of angular position for selected circles of pressure ports close to the jet orifice. Part (d) shows contours of constant pressure coefficient. These $\,C_p\,$ contour plots are useful as a convenient means for qualitative comparisons with the results of other test conditions or with the results of other experiments. Construction of the contour plots involved smoothing the data from the ray and circle plots.

From these figures some observations can be made about the effect on the flat-plate pressure distribution of changing the effective velocity ratio:

- 1. For all velocity ratios there is a positive pressure region upstream of the jet, but in general, the maximum pressure attained and the extent of this region decrease with increasing velocity ratio. This is probably due to the increased entrainment of crossflow fluid as the relative jet velocity increases.
- 2. The ray plots indicate a large radial component of the pressure gradient in the region close to and to the side of the jet. Recent analysis indicates that these gradients may be induced by the location of the counter-rotating vortices in the vicinity of the jet.
- 3. From the circle plots, it is seen that the angular position of the minimum pressure point moves upstream from about $110^{\rm O}$ for R=2.2 to about $75^{\rm O}$ for R=10. The value of the minimum pressure coefficient measured on the inner circle of ports is also a function of velocity ratio and appears to reach its lowest value of about -4.7 at the velocity ratio of 6.1. As a comparison, for potential flow over a circular cylinder the minimum pressure coefficient occurs at $90^{\rm O}$ and has a value of -1.5.
- 4. In the wake region along the $180^{\rm O}$ ray, the radial component of the pressure gradient is less than it is in the region to the side of the jet. The minimum pressure coefficient attained on this ray is about -2 for velocity ratios of 2 to 5. For velocity ratios above 5 the minimum value of C_p on this ray gradually becomes less negative with increasing velocity ratio, indicating a more rapid pressure recovery in the wake region for the larger velocity ratios. This is consistent with the physical observation that at the lower velocity ratios the deflected jet plume is closer to the flat plate, introducing a wake effect. At higher velocity ratios, the plume will be higher and the wake effect therefore less apparent.
- 5. An interesting behavior of the pressure distribution along the rays from $160^{\rm O}$ to $180^{\rm O}$ is noted for $R \ge 4$. There appears to be a small region close to and in the wake of the jet where a flat spot (very small radial component of the pressure gradient) or, perhaps even a hook, in the ray plot occurs. This behavior, however, is masked by the relatively large scatter of the data in the wake region.
- 6. An anomaly in the pressure distribution in the wake occurs for a velocity ratio of 3.9. The minimum pressures along the 170° and 180° rays, for example, are lower than the corresponding pressures for all other velocity ratios. The implication of this behavior is that there is a slower pressure recovery in the wake region for a velocity ratio of 4 than for the other velocity ratios.

7. For circles of pressure ports near the jet orifice, the value of C_p at a given angular position in the left half-plane $(-\theta)$ is consistently more negative than the value for the symmetry point in the right half-plane $(+\theta)$. This is due to misalinement between the nozzle and flat plate. The nozzle was displaced 0.02D in the negative Y-direction with respect to the point of intersection of the rays of pressure ports.

Most of the effects described above are also evident in the $\,C_p\,$ contour plots. For example, the slow pressure recovery in the wake region is described by contours that extend far downstream before intersecting the X-axis, whereas a more rapid pressure recovery is described by these same contours intersecting the X-axis much closer to the jet. It is also noted that the lobe of each contour appears to be swept back more for lower velocity ratios. Thus, the direction of maximum influence of the pressure field is a function of the velocity ratio.

The pressures at selected locations on the flat plate were also measured as a function of time for effective velocity ratios of 4.1 and 8. This was accomplished by sampling the output of the pressure transducer connected to a given port for approximately 10 seconds at a sampling rate of 31 times per second. These measurements were taken with the output of the pressure transducer filtered by a low pass filter and, also, unfiltered. For the latter case, the amplitude of pressure fluctuations with a frequency above about 15 cps should have been attenuated in the tubing connecting the static pressure port on the flat plate to the transducers. For the results of both cases, it is convenient to write the pressure coefficient at a given time as the sum of an average value over the time interval, of measurement and a fluctuation about this average value; $C_p = \overline{C}_p + C_p'$ where

 $\overline{C}_p = \frac{1}{n} \sum_{i=1}^n C_{p_i}$, and where n is the total number of times the pressure C_{p_i} at the port was measured. For this phase of the experiment, the jet-off pressures were not adequately measured so the values of p are not corrected for the jet-off condition but are instead referenced to free-stream pressure. The standard deviation of the fluctuations of the pressure coefficient about its mean value is also calculated:

$$\sigma = \left[\frac{1}{n} \sum_{i=1}^{n} \left(C_{p_i}'\right)^2\right]$$

Results of these calculations are presented in table 2 and in figures 12 to 15. Figure 12 shows the time variation of the pressure coefficient for selected pressure ports for an effective velocity ratio of 4.1 and for the output of the pressure transducer filtered.

TABLE 2.- PRESSURE-COEFFICIENT STANDARD DEVIATIONS

	Γ .		Pressu	re-coef	ficient	standar	d deviat	ions at	θ of	-	
r/D	200				120 ⁰	·——	r .	· ·	·	180 ^O	1700
	60 ⁰	75 ⁰					150 ⁰	160 ⁰	170 ⁰	180	-170 ⁰
		R = 4	.1; M _j	= 0.48;	pressu	re-tran	sducer	output f	iltered		
0.58	0.07	0.09	0.12	0.21	0.22	0.12	0.16	0.16	0.12	0.11	
.66		.07	.10	.17	.18	.12	.14	.14	.11	.09	0.12
.70		.05	.09	.14	.13	.09	.11	.11	.08	.07	.11
.80		.05	.06	.09	.09	.09	.11	.11	.09		.09
1.20		.02	.03	.05	.06	.06	.08	.09	.09		.12
2.50		.01	.01	.01	.02	.03	, .03	.05	.07	.06	.06
3.75		.00	.01	.01	.01	.02	.02	.03	.06	.09	.06
		R = 4.1	i; M _j =	= 0.48; p	ressur	e-trans	ducer o	utput un	filtered	l	
0.58	0.45	0.72	1.31	2.53	2.35	1.75	1.95	1.35	0.85	0.70	
.66		.60	1.12	1.96	2.02	1.54	1.57	1.25	.75	.59	0.78
.70		.58	1.04	1.72	1.94	1.54	1.60	1.26	.86	.73	.93
.80		.44	.82	1.33	1.70	1.40	1.29	1.17	.78		.81
1.20		.20	.37	.51	.73	.74	.71	.82	.70		
2.50		.04	.06	.07	.12	.19	.22	.34	.46		.49
3.75		.02	.03	.03	.04	.07	.10	.16	.30	.46	.43
		$\mathbf{R} = 8$.0; M _j	= 0.93;	pressu	re-tran	sducer	output f	iltered		
0.62	0.13	0.18	0.26	0.33	0.24	0.20	0.16	0.12	0.09	0.09	
.66		.12	.19	.25	.18	.15	.14	.10	.08	.07	0.09
.70		.11	.17	.22	.16	.14	.12	.09	.07	.07	.09
.80		.10	.15	.18	.15	.13	.12	.09	.07	.06	.08
1.00		.06	.08	.09	.09	.06	.07	.05	.04	.04	.05
2.00		.01	.02	.02	.02	.03	.04	.04	.03		
2.50		.01	.01	.01	.01	.01	.01	.02	.02	.03	
5.00		.01	.01	.01	.01	.01	.01	.01	.02	.02	
{		R = 8.0); Mj =	• 0.93; p	pressur	e-trans	ducer o	utput ur	filtered	ì	
1.00		0.53	0.74	0.92	0.65	0.55	0.59	0.48	0.36	0.25	0.24
2.00		.24	.33	.40	.40	.37	.53	.41	.27	.25	.31
3.50		.07	.13	.10	.14	.15	.20	.21	.22	.44	.46
5.00		.03	.08	.05	.07	.10	.12	.14	.14		.14

Results are shown for circles of pressure ports close to the jet orifice and for pressure ports on rays of $105^{\rm O}$ and $170^{\rm O}$. For the circles of ports close to the jet the maximum standard deviations occur at an angular location of $105^{\rm O}$, which coincides with the location of the point of minimum pressure (see fig. 6(c)). However, the standard deviation of the fluctuations decreases more rapidly with the distance along the $105^{\rm O}$ ray than along the $170^{\rm O}$ ray. Thus, for values of r/D greater than about 1, the maximum fluctuations occur in the wake region.

Results for the same test conditions, but with the low pass filters removed from the circuit, are shown in figure 13. The standard deviations of the fluctuations in the measured pressure are an order of magnitude greater, and higher frequencies are recorded, but the qualitative observations made for figure 12 are applicable. Results for R=8 are shown in figures 14 and 15. It appears that the fluctuations in C_p for R=8 are greater than those for R=4 in the region 60° to 105° but are less than those for R=4 in the wake region. These differences indicate that the distribution of turbulence on the plate is a function of effective velocity ratio.

The response of the pressure transducers connected to the static pressure ports on the jet exit ring indicate that the flow in the jet near the orifice had a relatively low level of turbulence as was expected. The fluctuations in the pressures measured on the jet exit ring were about three orders of magnitude smaller than those in the pressures measured on the plate near the orifice.

Although it was not the intent of this experiment to study the turbulent fluctuations of the pressure distributions on the flat plate, some feeling for the extent of the high turbulence region around the jet and a rough estimate of the turbulent intensities can be obtained from the values of σ listed in table 2. Contours of constant σ are plotted in figure 16 for velocity ratios of 4.1 and 8.0. Note that near the jet orifice, the shapes of these contours are similar to the shapes of the respective pressure contours. Since each data point for figures 4 to 11 represents a single random sample of the pressure-transducer output, the distribution of σ on the flat plate when the low pass filters are in the circuit is a measure of the distribution of uncertainty in the pressure coefficient. The large values of σ near the jet and in the wake help explain the scatter in the data for these regions.

While the C_p contour plots are the most convenient means of comparing one measured pressure distribution with another, care must be taken to recognize the effect of a given uncertainty in C_p on the location of each contour. If, for a given position on the flat plate, the pressure coefficient has an uncertainty of $\pm \Delta C_p$, then there is a corresponding uncertainty $\pm \Delta s$ in the location of the contour passing through that point. This uncertainty is given by $\Delta s = \Delta C_p/|grad C_p|$. Calculation of the uncertainty band width (2 Δs) for a contour thus requires a knowledge of the gradient of the pressure distribution at that point. From values of σ in table 2 and the measured pressure distribution for R = 3.9

from figure 6 the shaded uncertainty bands shown in figure 17(a) have been constructed. Note that the uncertainty bands become larger as the absolute magnitude of the pressure coefficient decreases; this is because the pressure gradient is decreasing with distance from the jet more rapidly than σ . Since no measurements of the pressure fluctuations were made for $\theta < 75^{\circ}$ no estimate of the uncertainty bands from fluctuations in C_p could be made in this region.

Another source of positional uncertainty in the $\,C_p\,$ contours would be inherent errors in the pressure transducers themselves. Although this type of error would be difficult to estimate accurately, a realistic value would be 0.5 percent of full-scale deflection. Such an error would result in the uncertainty bands shown in figure 17(b) for the same test conditions as those shown in figure 17(a). For this case of a constant uncertainty ΔC_p , the bandwidth at any location of the contour is determined solely by the magnitude of the gradient of $\,C_p\,$ at that location. The extreme width of the uncertainty bands for the contours -0.1, 0.0, and +0.1 except in the region near the jet indicates that using these contours to compare two separately determined pressure distributions could be misleading.

The C_p contours shown in figure 17 represent three different determinations of the pressure distribution at the same jet and free-stream Mach numbers for the last three test conditions listed in table 1. The differences appear not to be significant. Although not shown here, the differences in the pressure distributions for the first three conditions of table 1 indicated that the leading-edge extension was needed. It appears that the leading edge needs to be 8D to 10D upstream of the jet exit to obtain stable boundary-layer conditions in the region of interest. Therefore, all the data presented in this report was taken with the leading-edge extension added to the flat plate. The effects on the pressure distribution of changing the jet Mach number while holding the effective velocity ratio constant are shown in figure 18 for effective velocity ratios of 4, 5, 6, 7, 8, and 10. Using figure 17 as a guide for the uncertainties in contour locations, it is concluded that within the 0.45 to 0.95 range of jet Mach number studied this effect is small compared to the effect of changing the effective velocity ratio except in the wake of the jet. Note that the anomaly of relatively slow pressure recovery in the wake for R = 4.0 and $M_j = 0.75$ is supported by the results for R = 4.0 and $M_j = 0.44$.

Comparisons With the Results of Other Experiments

Comparisons of the results of the present experiment with those of several other experiments are provided in this section. These comparisons are made with $\,C_p\,$ contour plots in figures 19 to 23. Table 3 provides a summary of the test conditions for each experiment.

TABLE 3.- TEST CONDITIONS FOR VARIOUS EXPERIMENTS

**************************************	Jet diameter		Test-section	Flat-plate	Tun	nel speed	N _{Re}
Experiment	cm	in.	size	dimensions	m/sec	ft/sec	$(U_{\infty}D/\nu_{\infty})$
Soullier (ref. 21)	12.0	4.7	43D × 63D	21D × 21D	30 to 90	98.4 to 295	(24 to 72) × 10 ⁴
Present study	10.2	4.0	44D × 65D	$24D \times 27D$	23 to 62	75.5 to 203.4	(16 to 42) × 10 ⁴
Mosher (ref. 28)	5.1	2.0	54D (circular)	$24\mathrm{D} imes 33\mathrm{D}$	15	49	5.1 × 10 ⁴
Vogler (ref. 17)	2.5	.98	84D × 144D	24D × 40D	93	305	15 × 10 ⁴
Bradbury (ref. 18)	2.5	.98	102D × 138D	72D × 84D	18	59	3.0×10^4
Thompson (ref. 15)	2.5	.98	48D × 60D	Tunnel floor*	12	39	2.0×10^4
Ousterhout (ref. 23)	1.0	.39	75D × 107D	69D × 96D	28	91.9	1.8×10^{4}

^{*}Applied wind-tunnel boundary-layer suction.

The experiment chosen for the most extensive comparison was conducted by A. Soullier (ref. 21) at ONERA. As seen from table 3, the test conditions for this experiment are quite similar to those of the present study. Both effective velocity ratio and free-stream Mach number were matched for each comparison. Figure 19 presents the results. Although some of the low magnitude C_p contours (± 0.1 and ± 0.2) are as much as 1D apart, from the foregoing discussion it is believed that this behavior is not significant. Agreement between the contours for the larger values of C_p is fairly good, but sometimes there are significant though not drastic differences, for example, the differences shown in figure 19(a) for $C_p = -1.0$.

Figures 20 to 23 compare the results of the present study with those of the other experiments listed in table 3. For these comparisons effective velocity ratios were matched, but crossflow velocities were not matched. Agreement among the contour plots of Mosher (fig. 20), Thompson (fig. 21), Bradbury and Wood (fig. 21), Vogler (fig. 22), and the present study is fairly good, but for some contours there are discrepancies which would seem to exceed the experimental errors. For example, for R = 4 the contour indicates a more rapid pressure recovery in the wake for the results of Thompson and of Bradbury and Wood than for the results of Soullier, Mosher (ref. 28), and the present study. Agreement with the results of Ousterhout is not good (see fig. 23), particularly for the effective velocity ratio of 3.3. In this case, the results of Ousterhout represent a significantly more rapid pressure recovery in the wake and also a significant change in the shape of the contours to the side of the jet. The reasons for this are not presently known.

An order-of-magnitude range in both jet diameter and Reynolds number is seen among the experiments shown in table 3. The comparisons between the pressure distributions from the various experimenters indicate that the effect of jet diameter is small when distances are nondimensionalized with the jet diameter. Over the range of Reynolds number used in the experiments, this parameter also appears to have a small effect.

CONCLUSIONS

The results of a wind-tunnel investigation of the turbulent flow of a round subsonic jet of air exhausting perpendicularly through a flat plate into an essentially incompressible subsonic crossflow indicate the following conclusions about the static pressure distribution on the flat plate:

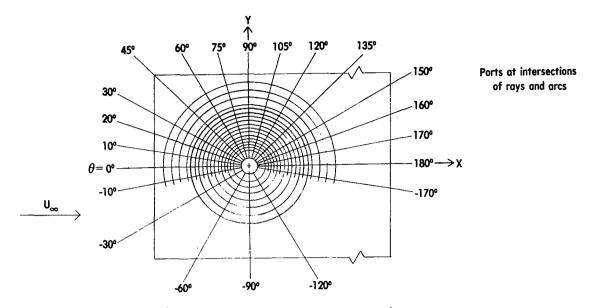
- 1. The dominant flow parameter for describing the pressure distribution on the plate is the effective velocity ratio. In particular, when distances are based on jet diameters, the pressure distribution does not appear to depend strongly on the jet size for jets up to 12 cm (5 in.) in diameter or on the jet Mach number. There do appear to be small but significant differences in the results of the experiments presented in this report. The parameters causing these secondary effects have not been identified.
- 2. There are regions of intense turbulence on the plate close to the jet orifice. For an effective velocity ratio of 4 the most severe turbulence is encountered at the point of minimum pressure. The intensity of the turbulence decays rapidly with distance to the side of the jet but persists farther into the wake region. The distribution of turbulence on the plate also appears to be a strong function of effective velocity ratio.
- 3. Some caution must be used when comparing flat-plate pressure distributions on the basis of contours of constant pressure coefficient. In regions of low pressure gradients, differences on the order of a jet diameter between contours are not unlikely even for the same test conditions.

Langley Research Center,
National Aeronautics and Space Administration,
Hampton, Va., March 31, 1975.

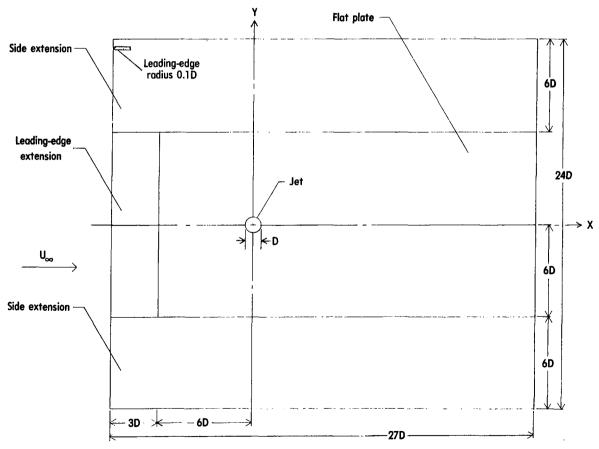
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(a) Pressure-port layout for $r/D \ge 1$.



(b) Original flat plate and extensions.

Figure 1.- Pressure-port locations and details of the flat plate.

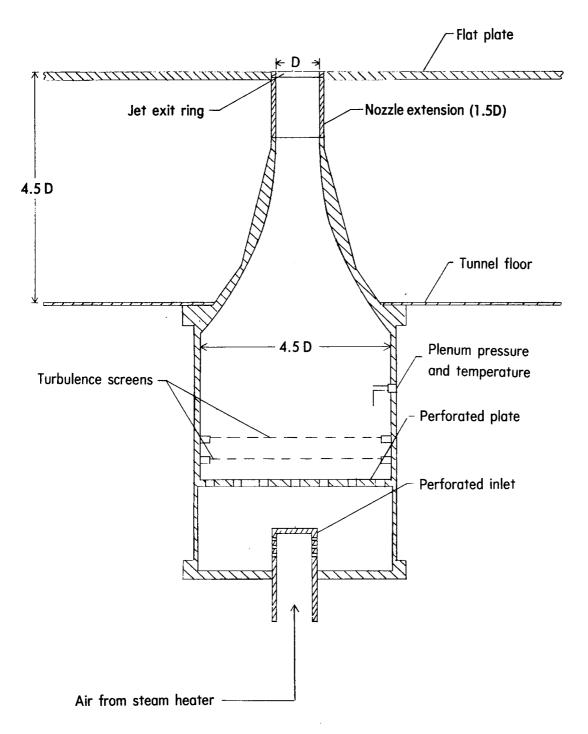
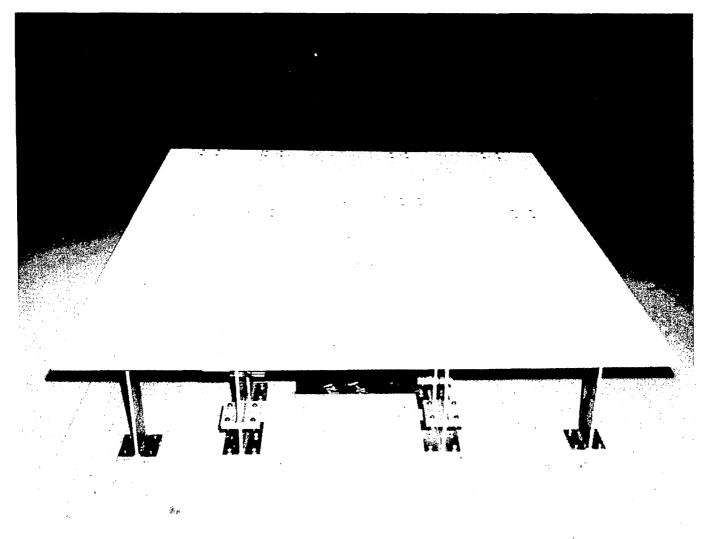


Figure 2.- Jet nozzle and plenum for plate 4.5D above tunnel floor. D = 10.16 cm (4.0 in.).



L-72-4893

Figure 3.- View of flat plate in V/STOL tunnel showing jet orifice.

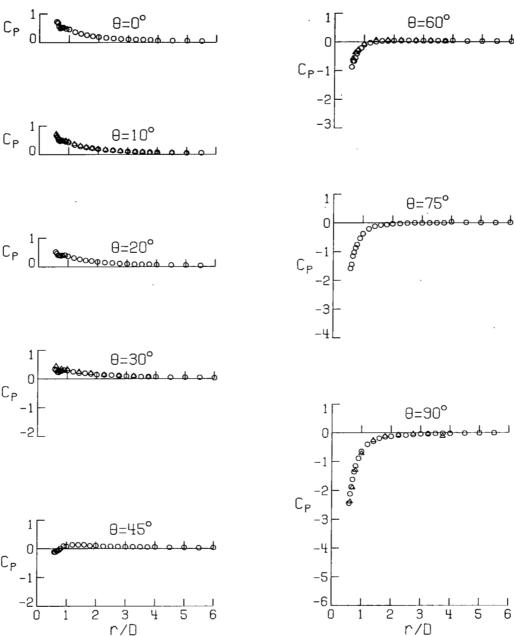
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105	-2.12	-2.51	-2.00	-1.91	-1.53	-1.50	-1.20	84	58	44	35	27	22
120	-2.54	-1.75	-1.46	-1.00	-1.40	-1.43	-1.05	-1.01	77	57	46	41	33
105	-1.39	-1.84	-1.77	-1.53	-1.17	-1.32	-1.22	86	83	71	72	47	48
150	-1.70	-1.40	-1.24	-1.15	-1.10	-1.16	43	83	96	73	68	53	55
Loo	-1.53	-1.39	-1.21	-1.50	-1.67	-1.25	92	83	83	72	05	58	48
170	-1.1+	-1.74	-1.55	-1.18	-1.15	-1.12	95	91	80	74	65	62	61
180	-l • 38	-1.72	-1.57	-1.33	-1.15	-1.13	-1.09	-1.08	94	87	74	66	53
-170	-1.93	-1.65	-1.74	-1.02	-1.42	-1.12	-1.27	-1.05	-1.04	91	70	78	56
-120		-3.92		-2.09		-2.10		-1.31		61		42	
-90		-2.39		-1.90		-1.29		72		27		13	
-6Ó		57		37		27		08		•06		•06	
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(a) C_p values.

Figure 4.- Pressure distribution on flat plate. R = 2.2; $M_j = 0.39$.

• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0)



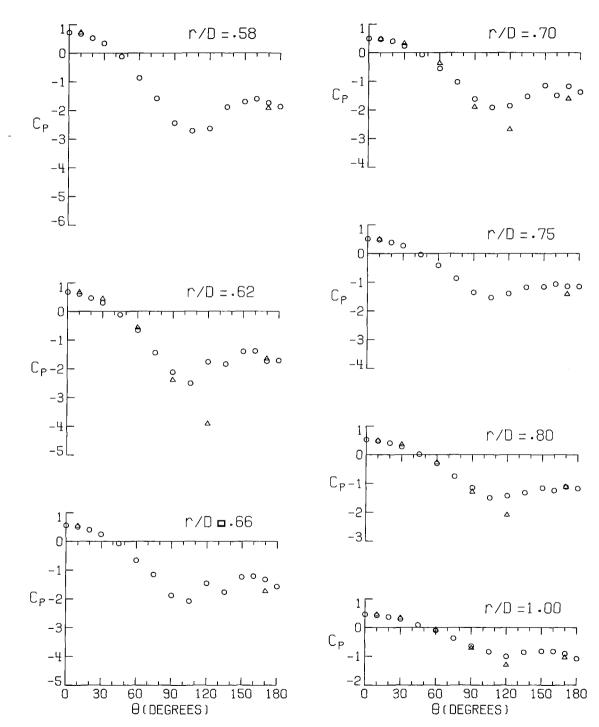
(b) Pressure distribution on rays.

Figure 4.- Continued.

• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0) 2 3 r/D r/D

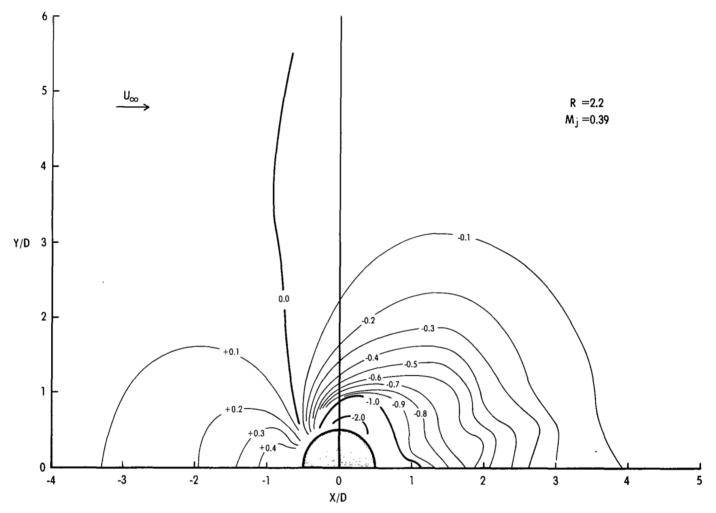
(b) Concluded.Figure 4.- Continued.

• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0)



(c) Pressure distribution on circles.

Figure 4.- Continued.



(d) Contours of constant C_p .

Figure 4.- Concluded.

r/D													
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23	• 44	• 35	• 30	.35	.34	•33	.32	•31	.24	•21	.18	.14	.14
33	.13	.10	.11	•15	•17	.17	.13	.20	-18	.15	•13	.13	.11
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90	-3.50	- i.l2	-2.72	-2.47	-2.04	-1.72	-1.35	-1.05	67	50	3/	30	24
105	-7.9+	-3.34	-3.17	-2.56	-2.58	-4.15	-1.60	-1.39	98	68	55	44	36
12j	ーチ・ラダ	-3.52	-3.23	-2.89	-2.74	-2.28	-2.11	-1.61	-1.11	92	71	57	53
135	- 3.29	-2.59	-2.66	-2.47	-2.27	-2.11	-1.57	-1.50	-1.29	99	86	70	57
150	-2.41	-2.41	-2.13	-1.01	-1.08	-1.79	-1.39	-1.04	-1.01	92	57	 74	55
150	-2.10	-1.9+	-1.68	-1.97	-1.53	-1.42	-1.23	43	31	7ó	71	58	56
170	-1.04	-1.33	-1.67	-1.50	د3.1-	-1.21	99	92	70	69	53	56	44
130	-1.80	-1.67	-1.68	-1.53	-1.47	-1.31	-1.13	93	80	04	68	52	38
-175	-1.37	-1.70	-1.67	-1.03	-1.43	-1.41	-1.22	-1.04	- • 34	58	67	45	38
-120		-4.4+		-3.46		-2.73		-1.79		98		65	
-90		-3.43		-2.64		-1.90		-1.07		52		30	
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0 13 20 30 45	.14 .12 .12 .10 .J4	.12 .12 .07 .08 .04 .04	.10 .10 .08 .08 .05	.08 .08 .07 .06 .03	.09 .08 .07 .06 .02	.07 .07 .06 .05 .02	.06 .05 .05 .04 .02	.04 .05 .05 .05 .01	.05 .04 .03 .03 .02	.04 .04 .04 .02 .01	.02 .04 .02 .03 .02	.00	
0 13 20 30 45 60 75	.14 .12 .12 .10 .34 32	21. 21. 20. 30. 80. 40.	.10 .10 .08 .08 .05 .00	.08 .08 .07 .06 .03 00	.09 .08 .07 .06 .02 01	.07 .07 .06 .05 .02 00	.06 .05 .05 .04 .02 .00	.04 .05 .05 .05 .01 00	.05 .04 .03 .03 .02 00	.04 .04 .04 .02 .01 01	.02 .04 .02 .03 .02 .00	.00	
0 13 20 30 45 60 75	.14 .12 .12 .10 .10 .14 02	.12 .12 .07 .08 .04 01 08	.10 .10 .08 .08 .05 .00 06	.08 .08 .07 .00 .03 -00 -05	.09 .08 .07 .06 .02 -01	.07 .07 .06 .05 .02 00 03	.06 .05 .05 .04 .02 .00 03	.04 .05 .05 .05 .01 00 03	.05 .04 .03 .03 .02 00	-04 -04 -04 -02 -01 -01 -32	.02 .04 .02 .03 .02 .00 00	.00 .00 .00 02	
0 13 20 30 45 60 75 90 105	.14 .12 .12 .10 .J4 J2 J9 19	.12 .12 .07 .08 .04 01 08 15	.10 .10 .08 .08 .05 .00 -06 -13	.08 .07 .00 .03 -00 -05 -11	.09 .08 .07 .06 .02 01 05 10	.07 .07 .06 .05 .02 00 03 04	.06 .05 .05 .04 .02 .00 -03 -01	.04 .05 .05 .01 00 03 06	. 05 . 04 . 03 . 03 . 02 - 00 - 03 - 05 - 08	.04 .04 .04 .02 .01 01 32 04	. 02 . 04 . 02 . 03 . 02 . 00 00 04 05	.00 .00 .00 02	
0 13 20 30 45 60 75 90 105 123	.14 .12 .12 .10 .J4 J2 J9 19 29	.12 .12 .07 .08 .04 01 08 15 23 33	.10 .10 .08 .08 .05 .00 06 13 19	.08 .08 .07 .00 .03 -00 -05 -11 -17	.04 .08 .07 .06 .02 -01 05 10	.07 .07 .06 .05 .02 00 03 04	.06 .05 .05 .04 .02 .00 03 01	.04 .05 .05 .01 00 03 06	. 05 . 04 . 03 . 03 . 02 00 03 05 08	.04 .04 .04 .02 .01 01 02 04	.02 .04 .02 .03 .02 .00 00 04 05	.00 .00 .00 02	05
0 13 20 30 45 60 75 90 105 123	.14 .12 .12 .10 .14 02 09 19 29 40	.12 .12 .07 .08 .04 01 08 15 23 33	.10 .10 .08 .08 .05 .00 -06 -13 -19 -28	.08 .08 .07 .06 .03 -00 -05 -11 -17 -23	.09 .08 .07 .06 .02 01 05 10	.07 .07 .06 .05 .02 00 03 04 13	.06 .05 .05 .04 .02 .00 03 01 11	.04 .05 .05 .01 00 03 06 10	.05 .04 .03 .02 00 03 05 08	.04 .04 .04 .02 .01 01 02 04 06 09	.02 .04 .02 .03 .02 .00 00 04 05	.00 .00 .00 02	
0 13 20 30 45 60 75 90 105 123 135	.14 .12 .12 .10 .J4 J2 J9 19 29 +J	.12 .12 .07 .08 .04 01 08 15 23 33	.10 .10 .08 .08 .05 .00 06 13 19 28 31	.08 .08 .07 .06 .03 -00 -05 -11 -17 -23 -29	.09 .08 .07 .06 .02 -01 05 10 15 19	.07 .07 .06 .05 .02 00 03 04 13 13	.00 .05 .05 .04 .02 .00 -03 -01 -11 -15	.04 .05 .05 .01 00 03 06 10 12	.05 .04 .03 .03 .02 -00 -03 05 08	.04 .04 .04 .02 .01 01 32 04 06 09 09	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 36 08	.00 .00 .00 02	-•05
0 13 20 30 45 60 75 90 105 123 135 153	-14 -12 -13 -10 -10 -12 -19 -19 -29 -40 -41	.12 .12 .07 .08 .04 01 08 15 23 33 41 41	.10 .10 .08 .08 .05 .00 -06 -13 -19 -28 -31 -33	.08 .08 .07 .00 .03 00 05 11 17 23 29	.09 .08 .07 .06 .02 01 05 10 15 19	.07 .06 .05 .02 00 03 04 13 13	.00 .05 .05 .04 .02 .00 -03 -01 -11 -15 -18	.04 .05 .05 .01 00 03 10 12 16	.05 .04 .03 .03 .02 -00 03 05 08 11	.04 .04 .04 .02 .01 01 02 04 06 09 09	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07	.00 .00 .00 02	-•05
0 13 20 30 45 60 75 90 105 123 135 153 153	-14 -12 -10 -10 -14 -12 -19 -29 -40 -47 -46 -32 -36	.12 .12 .07 .08 .04 01 23 33 41 41	.10 .10 .08 .08 .05 .00 06 13 19 28 31 33 23	.08 .08 .07 .00 .03 -005 -11 -17 -23 -29 -28	.09 .08 .07 .06 .02 01 05 10 15 19 24 22	.07 .07 .06 .05 .02 00 03 04 13 16 22	.00 .05 .05 .04 .02 .00 03 07 11 15 18	.04 .05 .05 .01 00 03 06 10 12 16	.05 .04 .03 .02 00 03 05 08 11 12 11	-04 -04 -04 -02 -01 -01 -04 -06 -09 -09 -08 -08	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07 07	.00 .00 .00 02 05 06	05
0 13 20 30 45 60 75 90 105 123 123 123 123 123 124 124 125 126 127 127 128	.14 .12 .10 .J4 J2 J9 J9 47 40 47 48 32	.12 .12 .07 .08 .04 01 23 33 41 41 35	.10 .10 .08 .05 .00 06 13 19 28 31 33 23	.08 .08 .07 .00 .03 -00 -05 -11 -17 -23 -29 -28 -20 -118	.09 .08 .07 .06 .02 -01 -05 -10 -15 -19 -24 -22 -16	.07 .07 .06 .05 .02 00 03 04 13 16 20 22 12	.06 .05 .05 .04 .02 .00 03 07 11 15 18 17 11	.04 .05 .05 .01 00 03 06 10 12 16 13	. 05 . 04 . 03 . 02 - 00 - 03 - 05 - 11 - 12 - 11 - 06 - 04 - 05	.04 .04 .04 .02 .01 01 02 04 09 09 09 08	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07	.00 .00 .00 02	-•05
0 13 20 30 45 60 75 90 105 123 135 153 163 170 180 -170	-14 -12 -10 -10 -14 -19 -19 -19 -19 -140 -140 -132 -136 -136 -136	.12 .12 .07 .08 .04 01 23 33 41 41	.10 .10 .08 .05 .00 06 13 19 28 31 33 26	.08 .08 .07 .00 .03 -005 -11 -17 -23 -29 -28	.09 .08 .07 .06 .02 -01 -105 -10 -15 -19 -24 -22 -16 -09	.07 .07 .06 .05 .02 00 03 04 13 16 22	.06 .05 .05 .04 .02 .00 03 07 11 15 18 17 11	.04 .05 .05 .01 00 03 06 10 12 16	.05 .04 .03 .02 00 03 05 08 11 12 11	-04 -04 -04 -02 -01 -01 -04 -06 -09 -09 -08 -08	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07 07	.00 .00 .00 02 05 06	05
0 13 20 30 45 60 75 90 105 123 135 153 160 170 160 -170 -120	-14 -12 -13 -10 -10 -19 -19 -29 -47 -46 -32 -38 -38	.12 .12 .07 .08 .04 01 23 33 41 41 35	.10 .10 .08 .08 .05 .00 06 13 19 28 31 33 23 23	.08 .08 .07 .00 .03 -00 -05 -11 -17 -23 -29 -28 -20 -118	.09 .08 .07 .06 .02 -01 -10 -15 -19 24 22 16 09 11	.07 .07 .06 .05 .02 00 03 04 13 16 20 22 12	.00 .05 .05 .04 .02 .00 -03 -01 -11 -15 -18 -17 -11 -10 -11	.04 .05 .05 .01 00 03 06 10 12 16 13	. 05 . 04 . 03 . 02 - 00 - 03 - 05 - 11 - 12 - 11 - 06 - 04 - 05	.04 .04 .04 .02 .01 01 02 04 09 09 09 08	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07 07	.00 .00 .00 02 05 06	05
0 13 20 30 45 60 75 90 105 123 153 153 160 170 160 -170 -120	-14 -12 -13 -10 -14 -17 -17 -29 -40 -47 -32 -36 -36 -37 -41 -17	.12 .12 .07 .08 .04 01 23 33 41 41 35	.10 .10 .08 .08 .05 .00 06 13 19 28 31 23 23 23 26 35 19	.08 .08 .07 .00 .03 -00 -05 -11 -17 -23 -29 -28 -20 -118	.09 .08 .07 .06 .02 01 15 19 24 22 16 09 11	.07 .07 .06 .05 .02 00 03 04 13 16 20 22 12	.00 .05 .05 .04 .02 .00 -03 -07 -11 -15 -18 -17 -11 -10 -11	.04 .05 .05 .01 00 03 06 10 12 16 13	. 05 . 04 . 03 . 02 - 00 - 03 - 05 - 11 - 12 - 11 - 06 - 04 - 05	.04 .04 .04 .02 .01 01 02 04 09 09 09 08	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07 07	.00 .00 .00 02 05 06	05
0 10 20 30 45 60 75 90 105 120 135 150 170 160 -170 -120 -90 -60	-14 -12 -10 -10 -19 -19 -29 -40 -47 -47 -32 -36 -38 -40 -19 -19	.12 .12 .07 .08 .04 01 23 33 41 41 35	.10 .10 .08 .08 .05 .00 06 13 19 28 31 33 23 24 35 19 29	.08 .08 .07 .00 .03 -00 -05 -11 -17 -23 -29 -28 -20 -118	.09 .08 .07 .06 .02 -01 05 -10 15 19 24 22 16 09 11	.07 .07 .06 .05 .02 00 03 04 13 16 20 22 12	.00 .05 .05 .04 .02 .00 03 07 11 15 18 17 11 10 11	.04 .05 .05 .01 00 03 06 10 12 16 13	. 05 . 04 . 03 . 02 - 00 - 03 - 05 - 11 - 12 - 11 - 06 - 04 - 05	.04 .04 .04 .02 .01 01 02 04 09 09 09 08	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07 07	.00 .00 .00 02 05 06	05
0 13 20 30 45 60 75 90 105 123 135 153 160 170 180 -170 -120	-14 -12 -13 -10 -14 -17 -17 -29 -40 -47 -32 -36 -36 -37 -41 -17	.12 .12 .07 .08 .04 01 23 33 41 41 35	.10 .10 .08 .08 .05 .00 06 13 19 28 31 23 23 23 26 35 19	.08 .08 .07 .00 .03 -00 -05 -11 -17 -23 -29 -28 -20 -118	.09 .08 .07 .06 .02 01 15 19 24 22 16 09 11	.07 .07 .06 .05 .02 00 03 04 13 16 20 22 12	.00 .05 .05 .04 .02 .00 -03 -07 -11 -15 -18 -17 -11 -10 -11	.04 .05 .05 .01 00 03 06 10 12 16 13	. 05 . 04 . 03 . 02 - 00 - 03 - 05 - 11 - 12 - 11 - 06 - 04 - 05	.04 .04 .04 .02 .01 01 02 04 09 09 09 08	. U2 . 04 . 02 . 03 . 02 . 00 00 04 05 06 08 07 07	.00 .00 .00 02 05 06	05

(a) C_p values.

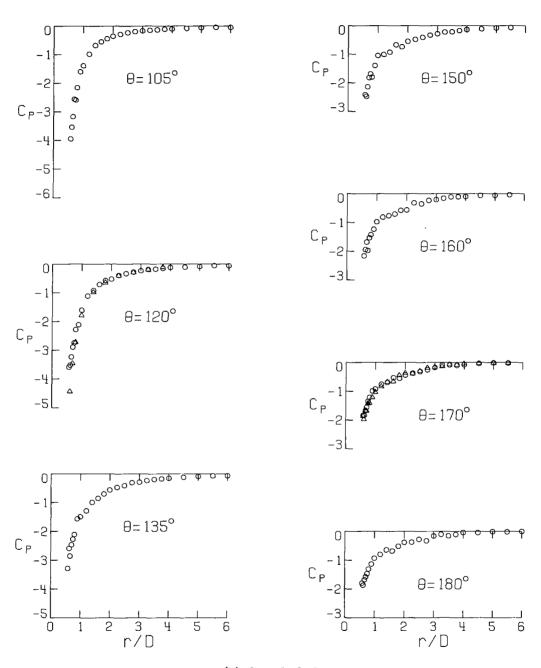
Figure 5.- Pressure distribution on flat plate. R = 2.8; $M_j = 0.44$.

RIGHT HALF+PLANE (0)△ LEFT HALF-PLANE (-0) 0 = 60° C_P-1 = 8

(b) Pressure distribution on rays.

Figure 5.- Continued.

RIGHT HALF-PLANE (0)△ LEFT HALF-PLANE (-0)



(b) Concluded.

Figure 5.- Continued.

• RIGHT HALF-PLANE (0) △ LEFT HALF-PLANE (-0) -5 -6L -5L

(c) Pressure distribution on circles.

30

0

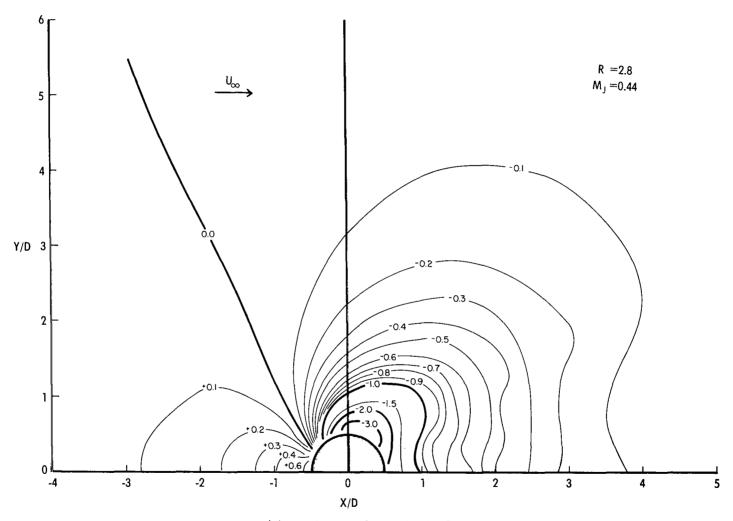
60

90

O(DEGREES)

Figure 5.- Continued.

0 (DEGREES)



(d) Contours of constant C_p .

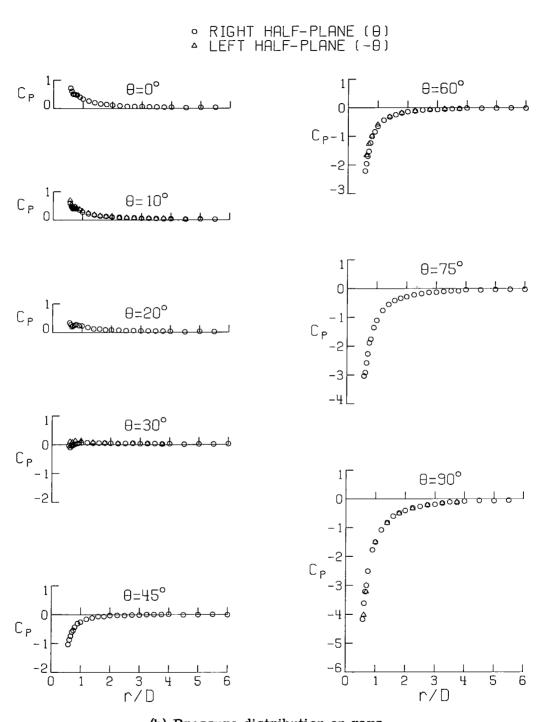
Figure 5.- Concluded.

θ r/D deg)	٤ر.	• 02	•06	.7u	• 75	• 30	.90	1.00	1.20	1.40	1.00	1.80	2.00
0	.71		.51	.49	• 4 9	.47	• 4 ů	. 33	• 26	•21	.16	-14	•12
10	• 59	• 47	•40	• 4 2	.43	•40	.35	-28	.22	.18	-14	.13	.11
23	د د .	.21	.20	.23	• 25	- 2d	. 24	.23	.13	.12	•12	• 09	•09
30	33	10	04	33	. J j	•02	دن.	.07	.07	• 05	.00	. 05	•05
45	-1.03	07	/4	60	54	43	31	26	10	11	06	06	03
აე	-2.21	-1.95	-1.09	-1.52	-1.23	-1.00	64	65	43	32	24	18	14
75	-3.33	-2.72	-2.50	-2.21	-1.33	-1.70	-1.35	-1.10	76	55	41	34	28
90	-4.17	-3.01	-3.21	-3.30	-2.51		-1.77	-1.50	-1.08	32	60	49	41
105	-4.32	- 3 . 3u	-3.44	-3.23	-2.44		-2.03	-1.79	-1.26	98	77	61	50
120	-3.54	-3.34	-3.17	-2.95	-2.75		-2.01	-1.85	-1.39	-1.12	84	73	60
135	44.ز-	-2. 1 d	-2.90	-2.40	-2.47		-1.95	-1.62	-1.24	96	76	65	56
150	-2.77	-2.54	-2.00	-2.31	-2.24	-2.00	-1.60	-1.47	-1.20	91	69	63	54
160	-2.09	-2.50	-2.49	-2.25	-2.23	-2.01	-1.00	-1.53	-1.17	-1.00	78	ó7	58
170	-2.20	-2.43	-2.48	-2.34	-2.22	-2.01	-1.75	-1.37	-1.19	88	86	86	61
180	-2.10	-2.23	-2.13	-2.34	-2.08	-1.99	-1.83	-1.46	-1.22	-1.22	94	82	77
-175	-1.98	-2.10	-4.21	~2.20	-2.38	-2.01	-1.77	-1.54	-1.31	-1.05	78	71	76
-120		-3.24		-2.05		-2.63		-1.07		-1.07		66	
-90		-4.02		-3.22		-2.03		-1.50		84		51	
-60		-1.00		-1.26		98		59		31		18	
- 30		.11		• Ü ö		.13		.12		.08		• 06	
-10	• 7 J	4ذ.	• 49	. 4 4	•47	.43	.3∂	.32	• 25	•19	.15	.12	•12

θ r/D (deg)	2.25	2.00	2.75	3.00	3.25	3.50	3.75	4.03	4.50	5.00	5.50	6.00	7.00
U	.10	• 07	. J7	.07	.06	.06	•05	.35	.03	•04	.04		
10	.10	.01	• 98	. 15	• 0 6	.06	.04	. 35	.03	.04	.03		
20	.08	•) 6	• 06	• 0 5	• 35	• U5	-04	•04	.02	.03	.02		
3 ს	.04	.04	• 35	.03	.04	.04	• U2	•03	.01	•02	.02	• 03	
45	32	03	51	31	• ၁ ၀	.01	.01	•02 `	01	00	•01	00	01
60	11	03	07	37	05	04	ذ0	01	02	02	01	02	
75	22	18	13	13	10	08	00	04	05	03	03	03	
90	31	27	21	19	15	11	13	08	0ó	07	05		
105	- • ÷ 3	30	29	24	27	18	15	13	06	08	07	07	
120	49	37	32	7.28	24	20	18	15	14	09	08	07	
135	40	37	32	30	22	19	19	13	13	09	08	08	05
150	44	- • 35	26	4	20	14	16	10	10	06	04		
160	40	37	33	22	23	15	15	09	09	03	05		
170	55	50	37	39	27	33	26	23	17	14	16		
180	74	55	47	50	42	43	43	31	35	30	33	30	12
-170	54	50	48	41	39	32	29	24	18	15			
-120	40		32		24		20						
-90	34		22		15		12						
-63	12		07		05		04						
-30	• 34		د٥.		•03		.02						
-10	• Oa	•04	• 07	•06	•06	.04	.06	.04	•02	•05			

(a) C_p values.

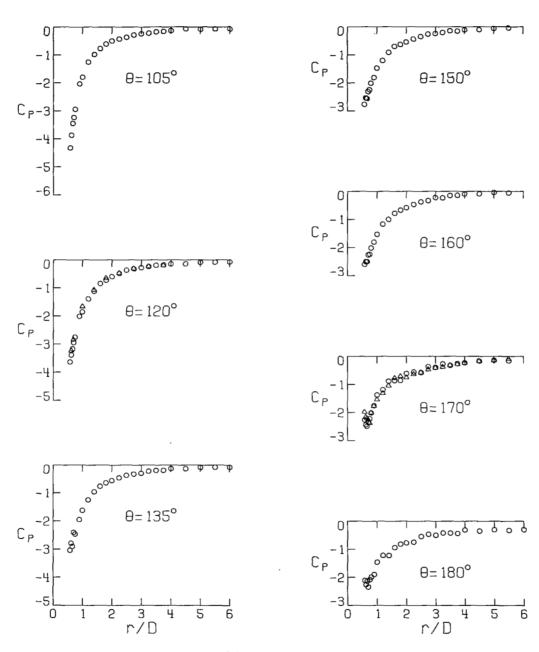
Figure 6.- Pressure distribution on flat plate. R = 3.9; $M_j = 0.75$.



(b) Pressure distribution on rays.

Figure 6.- Continued.

• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0)



(b) Concluded.

Figure 6.- Continued.

• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0) r/D=.58 r/D=.70 o 0 ۵ 0 -2 0 0 8 0 0 0 0 -5 -6L o r/D=.620 C_P-2 -4[-3 -5L -31 0 - 1 0 $C_{P}-2$ -3 -4 -51

(c) Pressure distribution on circles.

30

90

 Θ (DEGREES)

120

150

0

Figure 6.- Continued.

30

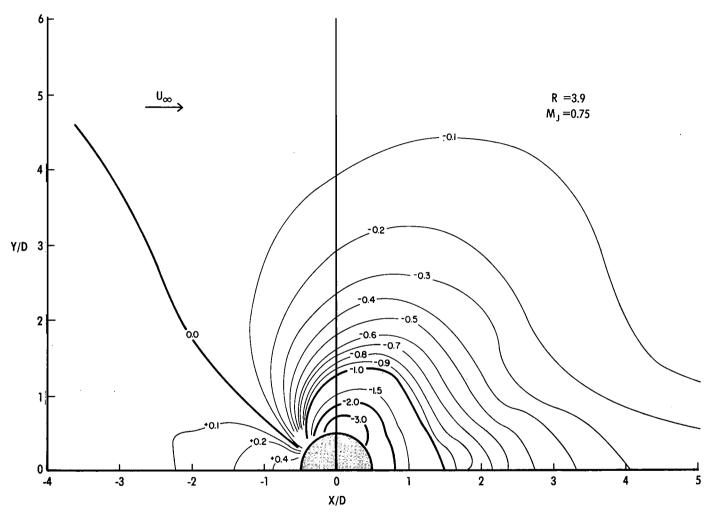
0

90

O(DEGREES)

120

150

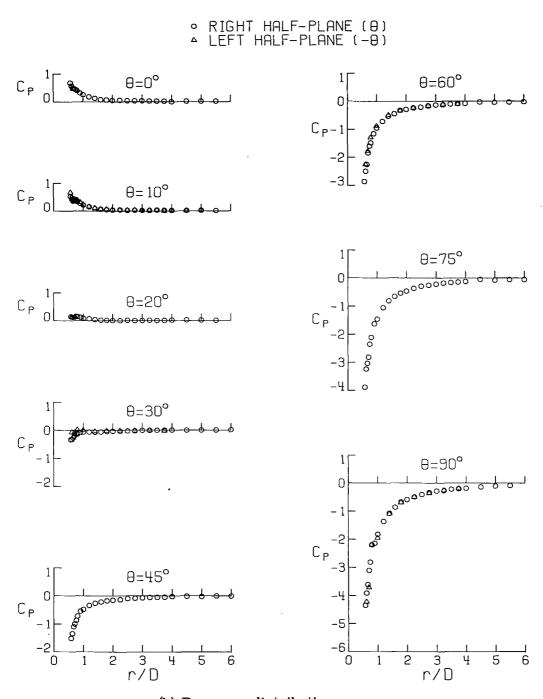


(d) Contours of constant C_p . Figure 6.- Concluded.

3	θ r/D (deg)	• 53	.62	• 06	.70	• 75	. 60	•90	1.00	1.20	1.40	1.60	1.80	2.00
20	J	7	.57	• 4 0	.47	•45	•40	• 32	. 26	-18	•13	•08	.07	.05
30	16	اد•	.42	• 36	.36	.37	.34	• 28	.22	- 15	.09	.07	• 05	-04
45	23	.10	د ا .	-13	-12	.16	- 15	.13	• 09	.07	• U 4	• 32	• 02	.01
50														
75	1 1													
90														
105														
120														
15	1 1													
150	, ,													
10	1 1													
170														1
1d0														
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$														
-12U -2.96 -2.03 -2.03 -2.20 -1.77 -1.0875 -1.0970 -6.0 -4.22 -3.71 -2.20 -1.73 -1.0970 -1.0970 -6.0 -2.23 -1.78 -1.3089513434300809 .01010503030809 .0101050303 -10 .36 .49 .43 .42 .42 .39 .31 .23 .16 .11 .07 .06 .0503 -10 .36 .49 .43 .42 .42 .39 .31 .23 .16 .11 .07 .06 .05030304 .03 .02 .02 .03 .02 .02 .03 .02 .02 .03 .02 .02 .03 .03 .02 .02 .03 .03 .02 .02 .03 .03 .02 .02 .03 .03 .02 .02 .03 .03 .02 .02 .03 .03 .03 .04 .03 .02 .02 .03 .03 .02 .02 .03 .03 .02 .02 .03 .03 .03 .02 .02 .03 .03 .02 .03 .03 .03 .02 .03 .03 .03 .03 .03 .03 .03 .03 .03 .03														
-90										• , ,		• > .		• • • •
-60	-90													1
-30	-60													
0 0 0 0 0 0 0 0 0 0	-3 0													
0	-10	• 56	• 49	. 43		• 42		.31		.16		.07		•05
10	<u> </u>													
10	1 0 1	2 • 20	2.50	2 • 75	3.30	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6.00	7.00
20	(deg)	٠.		·	<u>.</u>								6.00	7.00
30	(deg)	. 35	.03	•03	.04	•ú3	.02	• 02	•00	•02	•02	•02	6.00	7.00
45 13 10 09 08 06 05 02 02 00 01 03 60 24 21 17 13 13 10 08 08 03 04 03 02 75 37 29 25 22 18 16 14 12 05 08 06 05 90 49 40 33 28 27 22 20 18 11 09 105 55 49 40 36 30 25 25 20 16 13 11 09 120 54 51 38 35 29 28 24 21 17 14 12 10 135 49 39 35 30 20 20 18 19 16 11 10 00 </th <th>(deg)</th> <th>• 05 • 04</th> <th>.03</th> <th>•03 •03</th> <th>.04 .02</th> <th>•03 •02</th> <th>.02 .03</th> <th>.02</th> <th>•00 •02</th> <th>•02 •02</th> <th>.02</th> <th>•02 •02</th> <th>6.00</th> <th>7.00</th>	(deg)	• 05 • 04	.03	•03 •03	.04 .02	•03 •02	.02 .03	.02	•00 •02	•02 •02	.02	•02 •02	6.00	7.00
60	(deg) 3 10 20	.05 .04 .01	.03 .02 .01	• 03 • 03 • 00	•04 •02 •01	.03 .02 .02	.02 .03 .01	.02 .03 .01	.00 .02 .03	.02 .02 .03	.02 .03 .02	.02 .02 .01	<u>-</u>	7.00
75	(deg) 30	.05 .04 .01	.03 .02 .01	.03 .03 .00 03	.04 .02 .01	.03 .02 .02 01	.02 .03 .01 00	.02 .03 .01	.00 .02 .03	.02 .02 .03 00	.02 .03 .02	.02 .02 .01	•01	
105	(deg) 0 10 20 30 45	.05 .04 .01 04	.03 .02 .01 03	.03 .03 .00 03 09	.04 .02 .01 01	.03 .02 .02 01	.02 .03 .01 00	.02 .03 .01 .00	.00 .02 .03 .01	.02 .02 .03 00	.02 .03 .02 .00 02	.02 .02 .01 .00	•01 -•01	
120	(deg) 10 20 30 45 60	.05 .04 .01 04 13	.03 .02 .01 03 10	.03 .03 .00 03 09	.04 .02 .01 01 08	.03 .02 .02 01 06	.02 .03 .01 00 05 10	.02 .03 .01 .00 05	.00 .02 .03 .01 02 08	.02 .02 .03 00 02	.02 .03 .02 .00 02	.02 .02 .01 .00 00	.01 01 02	
135	(deg) 10 20 30 45 60 75	.05 .04 .01 04 13 24	.03 .02 .01 03 10 21 29	.03 .03 .00 03 09 17	.04 .02 .01 01 08 13	.03 .02 .02 01 06 13	.02 .03 .01 00 05 10 16	.02 .03 .01 .00 05 08	.00 .02 .03 .01 02 08	.02 .02 .03 00 02 03	.02 .03 .02 .00 02 04 08	.02 .02 .01 .00 00 03 06	.01 01 02	
150	(deg) 10 20 30 45 60 75 90 105	.05 .04 .01 04 13 24 37 49	.03 .02 .01 03 10 21 29 40	.03 .03 .00 03 09 17 25 33	.04 .02 .01 01 08 13 22	.03 .02 .02 01 06 13 18 27	.02 .03 .01 00 05 10 16	.02 .03 .01 .00 05 08 14	.00 .02 .03 .01 02 03 12	.02 .02 .03 00 02 03 05	.02 .03 .02 .00 02 04 08	.02 .02 .01 .00 00 03 06	•01 -•01 -•02 -•05	
100 30 29 24 20 15 14 12 11 06 05 04 170 30 39 38 21 25 23 17 22 18 14 11 180 40 39 29 40 31 26 20 19 18 23 24 14 08 -170 40 34 31 23 27 30 14 14 09 09 -120 50 36 26 19 -90 50 36 26 19 -60 24 17 13 10 -30 02 02 01 03	(deg) 10 20 30 45 60 75 90 105 120	.05 .04 .01 04 13 24 37 49 53	.03 .02 .01 03 10 21 29 40	.03 .03 .00 03 17 25 33	.04 .02 .01 01 08 13 22 28	.03 .02 .02 01 06 13 18 27 30	.02 .03 .01 00 05 10 16 22	.02 .03 .01 .00 05 08 14 20	.00 .02 .03 .01 02 08 12 18	.02 .02 .03 00 02 03 05 13	.02 .03 .02 .00 02 04 08 11	.02 .02 .01 .00 00 03 06 39	• 01 • 01 • 02 • 05 • 09	
170	(deg) 10 20 30 45 60 75 90 105 120 135	.05 .04 .01 04 13 24 37 49	.03 .02 .01 03 10 21 29 40 49	.03 .03 .00 -03 09 17 25 33 40 38	.04 .02 .01 01 08 13 22 26 36	.03 .02 .02 01 06 13 18 27 30	.02 .03 .01 00 05 10 16 22 25 28	.02 .03 .01 .00 05 08 14 20 25 24	.00 .02 .03 .01 02 08 12 18 20 21	.02 .02 .03 00 02 03 05 13 16	.02 .03 .02 .00 02 04 08 11	.02 .02 .01 .00 00 03 06 09 11	•01 •01 •02 •05 -09	03
180	(deg) 10 20 30 45 60 75 90 105 120 135	.05 .04 .01 04 13 24 37 49 53 54 49	.03 .02 .01 -03 10 21 29 40 49 51	.03 .03 .00 -03 09 17 25 33 40 38 35	.04 .02 .01 -01 -08 -13 -22 -26 -36 -35 -35	.03 .02 .02 01 06 13 18 27 30 29 29	.02 .03 .01 00 05 10 16 22 25 28 20	.02 .03 .01 .00 05 14 20 25 24 18	.00 .02 .03 .01 02 08 12 18 20 21 21	.02 .03 00 02 03 05 13 16 17	.02 .03 .02 .00 02 04 08 11 13	.02 .02 .01 .00 00 03 06 09 11 12 10	•01 •01 •02 •05 -09	03
-170	(deg) 10 20 30 45 60 75 90 105 120 135	.05 .04 .01 04 13 24 37 49 53 54 49	.03 .02 .01 03 10 21 29 40 49 51 39 39	.03 .03 .00 03 17 25 33 40 38 35	.04 .02 .01 01 08 13 22 26 36 35 30 19	.03 .02 .02 01 06 13 18 27 30 29 20 16	.02 .03 .01 00 05 10 16 22 25 28 20 14	.02 .03 .01 .00 05 14 20 25 24 18	.00 .02 .03 .01 02 08 12 18 20 21	.02 .03 00 02 03 05 13 16 17	.02 .03 .02 .00 02 04 08 11 13 14 11	.02 .02 .01 .00 00 03 06 09 11 12 10	•01 •01 •02 •05 -09	03
-120504432252519362619103002020103	(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170	. 05 . 04 . 01 04 13 24 37 53 54 49 49 49 49 39 39 30	.03 .02 .01 03 10 21 29 40 49 51 39 32	.03 .03 .00 03 09 17 25 33 40 38 35 23 23	.04 .02 .01 01 08 13 22 28 36 35 30 19	.03 .02 .02 01 06 13 27 30 29 26 15	.02 .03 .01 00 05 10 16 22 25 28 20 14	.02 .03 .01 .00 05 14 20 25 24 18	.00 .02 .03 .01 02 08 12 19 20 21 09 11	.02 .02 .03 00 02 03 05 13 16 17 16	.02 .03 .02 .00 02 04 11 13 14 11	.02 .02 .01 .00 00 03 06 39 11 12 10	-01 -01 -02 -05 -09 -10 -08	03 06
-90 50 36 26 19 -60 24 17 13 10 -30 02 02 01 03	(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180	. 05 . 04 . 01 04 13 24 37 49 53 34 49 39 30 30 30 40	.03 .02 .01 03 10 21 29 49 51 39 39 39	.03 .03 .00 -03 17 25 33 40 38 23 24 35	.04 .02 .01 01 08 13 22 36 35 30 19 20 21	.03 .02 .02 01 06 13 18 27 30 29 20 15 25	.02 .03 .01 00 05 10 16 22 25 28 20 14 14	.02 .03 .01 .00 05 08 14 20 25 24 18 12 17	.00 .02 .03 .01 02 08 12 18 20 21 19 09	.02 .02 .03 00 02 03 05 13 16 17 16 07	.02 .03 .02 .00 02 04 08 11 13 14 11 06 05	.02 .02 .01 .00 00 03 06 39 11 12 10	-01 -01 -02 -05 -09 -10 -08	03 06
-6024171310030103	(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170	.05 .04 .01 04 13 24 37 53 54 49 30 30 40	.03 .02 .01 03 10 21 29 49 51 39 39 39	.03 .03 .00 -03 09 17 25 33 40 38 23 24 35	.04 .02 .01 01 08 13 22 36 35 30 19 20 21	.03 .02 .02 01 06 13 18 27 30 29 20 16 15 25 25	.02 .03 .01 00 05 10 16 22 25 28 20 14 14	.02 .03 .01 .00 05 08 14 20 25 24 18 12 12 17	.00 .02 .03 .01 02 08 12 18 20 21 19 09	.02 .02 .03 00 02 03 05 13 16 17 16 07	.02 .03 .02 .00 02 04 08 11 13 14 11 06 05	.02 .02 .01 .00 00 03 06 39 11 12 10	-01 -01 -02 -05 -09 -10 -08	03 06
-30 02 01 03	(deg) 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120	.05 .04 .01 04 13 24 37 49 53 54 30 30 30 30	.03 .02 .01 03 10 21 29 49 51 39 39 39	.03 .03 .00 03 17 25 33 40 38 23 24 38	.04 .02 .01 01 08 13 22 36 35 30 19 20 21	.03 .02 .02 01 06 13 18 27 30 29 16 15 25 31	.02 .03 .01 00 05 10 16 22 25 28 20 14 14	.02 .03 .01 .00 05 14 20 25 24 18 12 17 20	.00 .02 .03 .01 02 08 12 18 20 21 19 09	.02 .02 .03 00 02 03 05 13 16 17 16 07	.02 .03 .02 .00 02 04 08 11 13 14 11 06 05	.02 .02 .01 .00 00 03 06 39 11 12 10	-01 -01 -02 -05 -09 -10 -08	03 06
	(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120 -90	.05 .04 .01 04 13 24 37 49 53 54 49 30 30 40 40 40	.03 .02 .01 03 10 21 29 49 51 39 39 39	.03 .03 .00 03 17 25 33 40 38 23 24 38 24 38	.04 .02 .01 01 08 13 22 36 35 30 19 20 21	.03 .02 .02 01 06 13 18 27 30 29 20 16 15 25 31	.02 .03 .01 00 05 10 16 22 25 28 20 14 14	.02 .03 .01 .00 05 14 20 25 24 18 12 17 20	.00 .02 .03 .01 02 08 12 18 20 21 19 09	.02 .02 .03 00 02 03 05 13 16 17 16 07	.02 .03 .02 .00 02 04 08 11 13 14 11 06 05	.02 .02 .01 .00 00 03 06 39 11 12 10	-01 -01 -02 -05 -09 -10 -08	03
-LJ -JJ -JJ -OJ -OJ -OJ -OJ -OJ -OJ -OJ	(deg) 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120 -90 -60	. 05 . 04 . 01 04 13 24 37 53 54 49 30 30 40 50 50 50 50	.03 .02 .01 03 10 21 29 49 51 39 39 39	.03 .03 .00 -03 09 17 25 33 40 38 23 23 24 38 24 38	.04 .02 .01 01 08 13 22 36 35 30 19 20 21	.03 .02 .02 01 06 13 27 30 29 26 15 25 31 27	.02 .03 .01 00 05 10 16 22 25 28 20 14 14	.02 .03 .01 .00 05 08 14 20 25 24 18 12 17 20 14	.00 .02 .03 .01 02 08 12 18 20 21 19 09	.02 .02 .03 00 02 03 05 13 16 17 16 07	.02 .03 .02 .00 02 04 08 11 13 14 11 06 05	.02 .02 .01 .00 00 03 06 39 11 12 10	-01 -01 -02 -05 -09 -10 -08	03

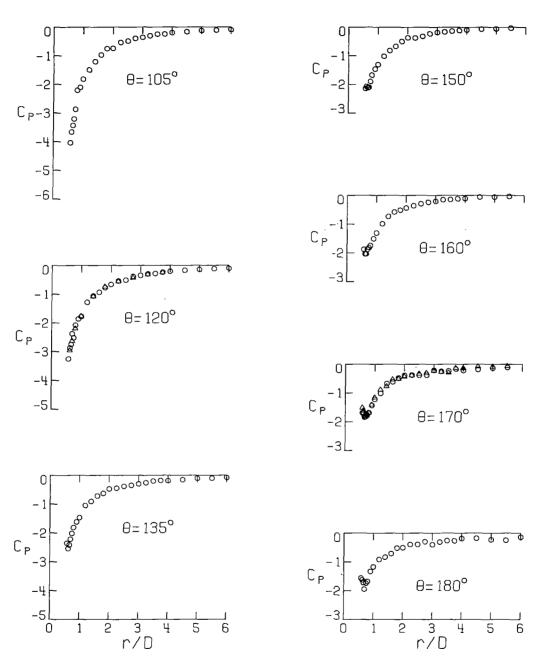
(a) C_p values.

Figure 7.- Pressure distribution on flat plate. R = 5.1; $M_j = 0.95$.



(b) Pressure distribution on rays.

Figure 7.- Continued.



(b) Concluded.

Figure 7.- Continued.

(c) Pressure distribution on circles.

0

30

90

(DEGREES)

120

-5L

30

90

O(DEGREES)

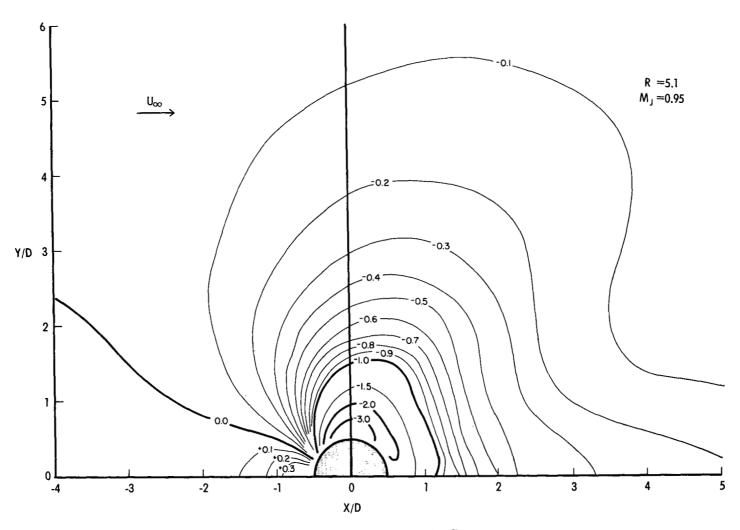
120

60

150

180

Figure 7.- Continued.



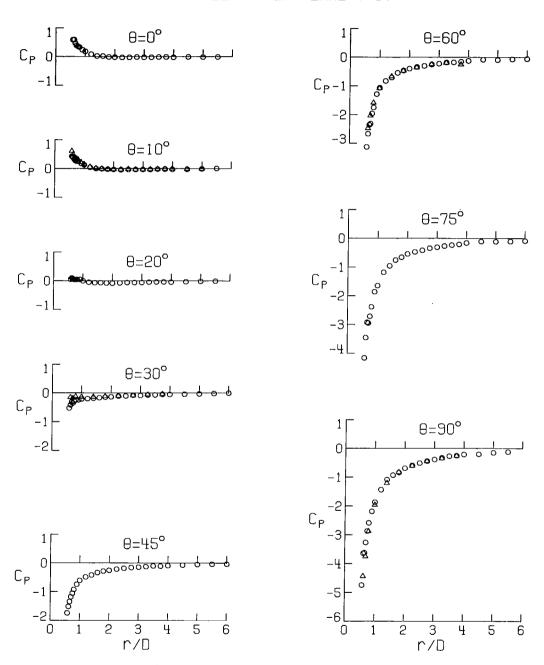
(d) Contours of constant C_p .

Figure 7.- Concluded.

θ r/D (deg)	• 58	•62	•60	.70	.75	•80	•90	1.00	1.20	1.40	1.60	1.80	2.00
	•59	.54	. 47	9 ز .	. 3 ó	•33	. 24	-18	.08	•02	•02	01	02
10	ذُهُ.	.34	.33	.28	•25	• 26	.18	.13	•04	.00	01	02	03
20	ون. دن.	.10	.07	• 35	.04	•05	.05	00	05	07	07	08	08
30	→. 52	42	32	35	~.29	4	24	21	20	19	17	15	13
45	-1.74	-1.51	-1.34	-1.18	-1.05	92	74	oi	48	42	33	28	25
60	-3.13	-2.67	-2.36	-2.31	-1.90	-1.75	-1.28	-1.06	82	71	54	46	39
75	-4.10	-3.45	-2.92	-2.95	-2.71	-2.38	-1.86	-1.65	-1.18	95	75	65	54
90	-4.74	-3.64	-3.62	-3.20	-2.87	-2.58	-2.18	-1.80	-1.43	-1.08	93	83	69
105	-3.83	-3.60	-3.23	-3.00	-2.53	-2.42	-2.25	-1.83	-1.47	-1.16	-1.05	89	75
120	-2.65	-2.59	-2.52	-2.52	-2.13	-2.19	-1.81	-1.52	-1.27	-1.09	49	74	72
135	-2.40	-2.00	-1.97	-1.98	-1.97	-1.69	-1.52	-1.17	95	82	68	56	52
150	-1.86	-1.93	-1.76	-1.38	-1.63	-1.58	-1.32	-1.22	42	68	57	44	38
160	-1.53	-1.71	-1.78	-1.82	-1.58	-1.49	-1.22	-1.04	73	65	54	40	37
170	-1.43	-1.57	-1.64	-1.51	-1.57	-1.45	-1.29	-1.03	83	64	43	42	42
180	-1.35	-1.39	-1.61	-1:57	-1.52	-1.47	-1.32	-1.09	74	67	49	38	39
-170	-1.46	-1.51	-1.59	-1.62	-1.50	-1.48	-1.37	-1.03	86	59	53	50	37
-120		-2.92		-2.10		-2.20		-1.66		-1.17	-	86	
-90		-4.44		-3.75		-2.88		-1.90		-1.21		83	
-60		-2.48		-2.04		-1.59		-1.09		67		47	
-30		15		18		13		14		14		12	
-10	•61	•48	.36	. 30	.34	•30	.21	.15	.05	.01	01	02	03
LL													
N ~/nl													
θ r/D (deg)	2.25	2.50	2.75	3.00	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6,00	7.00
(deg)	2.25	2.50	2.75	3.00	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6,00	7.00
			<u></u>						4.50 01		. ,	6,00	7.00
(deg)	03 04	2.50 03 03	2.75 02 03	3.00 J2 J3	3.25 02 03	3.50 01 02	3.75 01 02	4.00 02 02		5.00 01 01	5.50 01	6.00	7.00
(deg)	03	03	02	02	02	01	01	02	01	01	01	6,00	7.00
(deg)	03 04	03 03	02 03	J2 J3	02 03	01 02	01 02	02 02	01 02	01 01	01 .01	6.00	7.00
(deg) 0 10 20	03 04 07	03 03 00	02 03 06	J2 U3 O6	02 ⁻ 03 04	01 02 04	01 02 03	02 02 04	01 02 03	01 01 01	01 .01 01		7.00
(deg) 0 10 20 30	03 04 07 11	03 03 00 10	02 03 06 08	J2 03 06 09	02 03 04 07	01 02 04 06	01 02 03 05	02 02 04 05	01 02 03 04	01 01 01 03	01 .01 01 01	•00	·
(deg) 0 10 20 30 45	03 04 07 11 21	03 03 00 10 18	02 03 06 08 16	J2 03 06 09 14	02 03 04 07 12	01 02 04 06 10	01 02 03 05 11	02 02 04 05 09	01 02 03 04 07	01 01 01 03 04	01 .01 01 01	-00 03	·
(deg) 10 20 30 45 60	03 04 07 11 21	03 03 00 10 13 30	02 03 06 08 16 25	J2 03 06 09 14	02' 03 04 07 12 18	01 02 04 06 10 17	01 02 03 05 11 15	02 02 04 05 09 12	01 02 03 04 07	01 01 01 03 04 09	01 .01 01 01 03 06	-00 03 05	·
(deg) 10 20 30 45 60 75	03 04 07 11 21 34	03 03 00 10 18 30 41	02 03 06 08 16 25 34	J2 03 06 09 14 22 30	02' 03 04 07 12 18	01 02 04 06 10 17 23	01 02 03 05 11 15	02 02 04 05 09 12 16	01 02 03 04 07 08	01 01 01 03 04 09	01 01 01 01 03 06	-00 03 05	·
(deg) 0 10 20 30 45 60 75 90	03 04 07 11 21 34 47	03 03 00 10 13 30 41	02 03 06 08 16 25 34	U2 U3 06 09 14 22 3U 38	02 03 04 07 12 18 26	01 02 04 06 10 17 23	01 02 03 05 11 15 20	02 02 04 05 09 12 16	01 02 03 04 07 08 10	01 01 01 03 04 09 11	01 -01 01 03 06 10	- 00 03 05 08	04
(deg) 10 20 30 45 60 72 90 105	03 04 07 11 21 34 47 59	03 03 00 10 13 30 41 51	02 03 06 08 16 25 34 44	J2 U3 06 09 14 22 3U 38 42	02' 03 04 07 12 18 26 32	01 02 04 06 10 17 23 27	01 02 03 05 11 15 20 25	02 02 04 05 09 12 16 21	01 02 03 04 07 08 10 20	01 01 01 03 04 09 11	01 .01 01 03 06 10 12	-00 -03 -05 -08	·
(deg) 0 10 20 30 45 60 75 90 105 120	03 04 07 11 21 34 47 59 03	03 03 00 10 13 30 41 51	02 03 06 08 16 25 34 44 47	J2 03 06 09 14 22 30 38 42 38	02 ⁻ 0304071218263237	01 02 04 06 10 17 23 27 32	01 02 03 05 11 15 20 25 31	02 02 04 05 09 12 16 21 25	01 02 03 04 07 08 10 20 22	01 01 03 04 09 11 15	01 .01 01 03 06 10 12 15	-00 03 05 08 12	04
(deg) 0 10 20 30 45 60 75 90 105 120 135	03 04 07 11 21 34 47 59 03	03 03 00 10 13 30 41 51 53 55	02 03 06 08 16 25 34 44 47	J2 03 06 09 14 22 30 38 42 38 25	02'030407121826323735	01 02 04 06 10 17 23 27 32 31	01 02 03 05 11 15 20 25 31 27	02 02 04 05 09 12 16 21 25 22	01 02 03 04 07 08 10 20 22	01 01 03 04 09 11 15 17	01 -01 -01 -03 -06 -10 -12 -15 -15	-00 03 05 08 12	04
(deg) 10 20 30 45 60 75 90 105 120 135 150	03 04 07 11 21 34 47 59 52 52	03 03 00 10 13 30 41 51 53 55	02 03 06 08 16 25 34 44 47 45	J2 U3 06 09 14 22 30 38 42 38 25 15	02'03040712182632373524	01 02 04 06 10 17 23 27 32 31 26	01 02 03 05 11 15 20 25 31 27 19	02 02 04 05 09 12 16 21 25 22	01 02 03 04 07 08 10 20 22 20	01 01 03 04 09 11 15 17 17	01 -01 01 03 06 10 12 15 15	-00 03 05 08 12	04
(deg) 10 20 30 45 60 75 90 105 120 135 150 160	03 04 07 11 21 34 59 03 52 4d 31	03 00 10 13 30 41 51 53 55 33	02 03 06 08 16 25 34 47 45 32 22	J2 J3 O6 O9 14 22 38 42 38 25 15 17	02'0304071218263237352415	01 02 04 06 10 17 23 27 32 31 26 11	01 02 03 05 11 15 20 25 31 27 19	02 04 05 09 12 16 21 25 22 21	01 02 03 04 07 08 10 20 22 20 14	01 01 03 04 09 11 15 17 17 13 06	01 01 01 03 06 10 12 15 15 10	-00 03 05 08 12	04
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170	03 04 07 11 21 34 47 59 03 52 48 31 30	03 00 10 13 30 41 51 53 55 33 25	02 03 06 08 16 25 34 44 47 45 32 22 17	J2 J3 O6 O9 14 22 30 38 42 38 25 15 17 19	02030407121826373524151318	01 02 04 06 10 17 23 27 32 31 26 11	01 02 03 05 11 15 20 25 31 27 19 08	02 04 05 09 12 16 21 25 22 21 10	01 02 03 04 07 08 10 20 22 20 14 08	01 01 03 04 09 11 15 17 17 13 06	01 01 01 03 06 10 12 15 15 10 04	-00 -03 -05 -08 -12 -12 -110	04
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180	03 04 07 11 21 34 47 59 03 52 48 31 30	03 03 00 10 18 30 41 51 53 55 33 25 21 24	02 03 06 08 16 25 34 44 47 45 32 22 17 19	J2 03 06 09 14 22 30 38 42 38 25 15 17 19 19	02'03040712182632373524151318	01 02 04 06 10 17 23 27 32 31 26 11	01 02 03 05 11 15 20 25 31 27 19 08 07	02 02 04 05 09 12 16 21 25 22 21 10 13	01 02 03 04 07 08 10 20 22 20 14 08	01 01 03 04 09 11 15 17 17 13 06 04	01 01 01 03 06 10 12 15 15 10 04	-00 -03 -05 -08 -12 -12 -110	04
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170	03 04 07 11 21 34 47 59 52 48 31 30 30	03 03 00 10 18 30 41 51 53 55 33 25 21 24	02 03 06 08 16 25 34 44 47 45 32 22 17 19 23 20	J2 03 06 09 14 22 30 38 42 38 25 15 17 19 19	02'0304071218263237352415131818	01 02 04 06 10 17 23 27 32 31 26 11	01 02 03 05 11 15 20 25 31 27 19 08 07	02 02 04 05 09 12 16 21 25 22 21 10 13	01 02 03 04 07 08 10 20 22 20 14 08	01 01 03 04 09 11 15 17 17 13 06 04	01 01 01 03 06 10 12 15 15 10 04	-00 -03 -05 -08 -12 -12 -110	04
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 -170 -120	03 04 07 11 21 34 59 53 52 30 30 32 36	03 03 00 10 18 30 41 51 53 55 33 25 21 24	02 03 06 08 16 25 34 44 47 45 32 22 17 19	J2 03 06 09 14 22 30 38 42 38 25 15 17 19 19	02'030407121826323735241513181814	01 02 04 06 10 17 23 27 32 31 26 11	01 02 03 05 11 15 20 25 31 27 19 08 07 20	02 02 04 05 09 12 16 21 25 22 21 10 13	01 02 03 04 07 08 10 20 22 20 14 08	01 01 03 04 09 11 15 17 17 13 06 04	01 01 01 03 06 10 12 15 15 10 04	-00 -03 -05 -08 -12 -12 -110	04
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120 -90 -60 -30	03 04 07 11 21 34 59 03 52 4d 31 30 32 36 36	03 03 00 10 18 30 41 51 53 55 33 25 21 24	02 03 06 08 16 25 34 47 45 32 22 17 19 23 20 45	J2 03 06 09 14 22 30 38 42 38 25 15 17 19 19	02030407121826373524151818181906	01020406101723273231261115121812	01020305111520253127190807202314262504	02 02 04 05 09 12 16 21 25 22 21 10 13 20	01020304070810202220140806101509	01 01 03 04 09 11 17 17 13 06 04	01 01 01 03 06 10 12 15 15 10 04	-00 -03 -05 -08 -12 -12 -11	04
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120 -90 -60	03 04 07 11 21 34 47 59 33 52 48 30 30 32 36 61	03 03 00 10 18 30 41 51 53 55 33 25 21 24	020306081625344447453222171923204545	J2 03 06 09 14 22 30 38 42 38 25 15 17 19 19	0203040712182632373518181814373519	01 02 04 06 10 17 23 27 32 31 26 11	01 02 03 05 11 15 20 25 31 27 19 08 07 20 23 14 26	02 02 04 05 09 12 16 21 25 22 21 10 13	01 02 03 04 07 08 10 20 22 20 14 08	01 01 03 04 09 11 15 17 17 13 06 04	01 01 01 03 06 10 12 15 15 10 04	-00 -03 -05 -08 -12 -12 -11	04

(a) C_p values.

Figure 8.- Pressure distribution on flat plate. R = 6.1; $M_j = 0.94$.

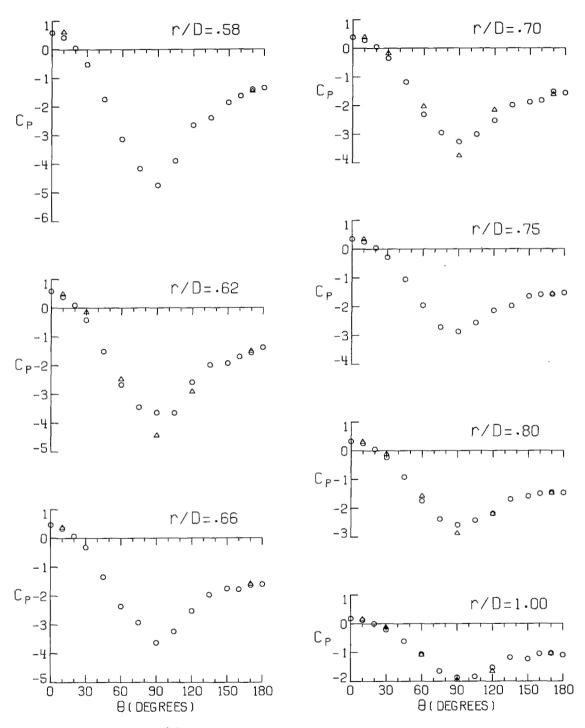


(b) Pressure distribution on rays.

Figure 8.- Continued.

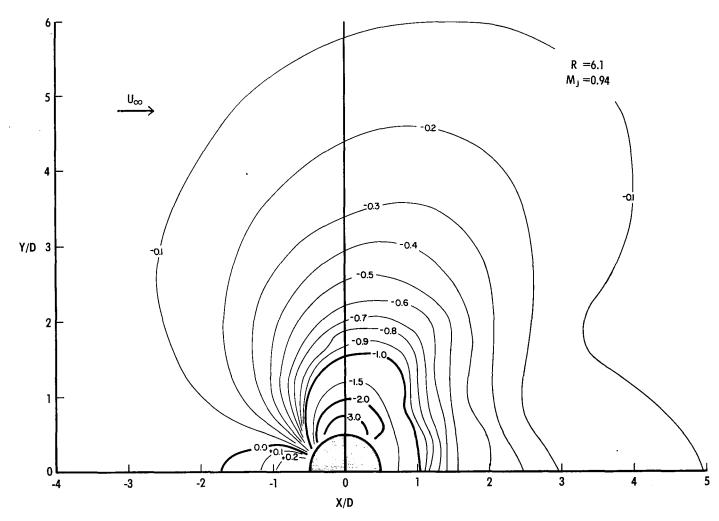
Figure 8.- Continued.

(b) Concluded.



(c) Pressure distribution on circles.

Figure 8.- Continued.



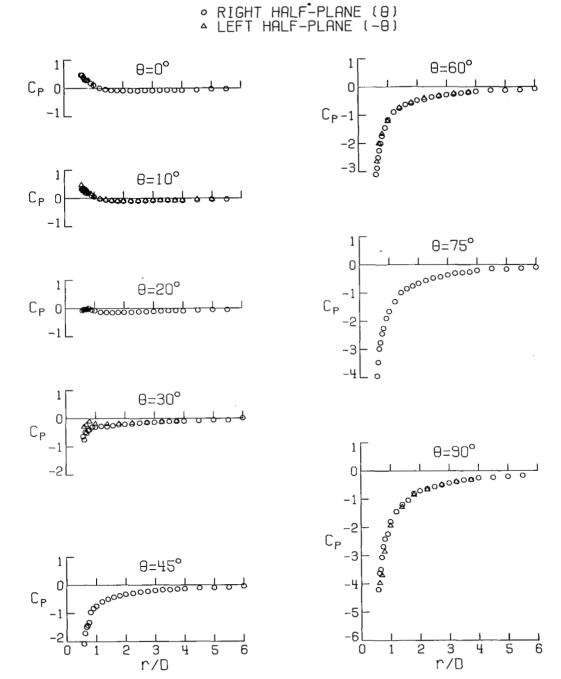
(d) Contours of constant C_p .

Figure 8.- Concluded.

e r/D	• 58	•62	. 66	.70	.75	.80	•90	1.00	1.20	1.40	1.60	1.80	2.00
(deg)	• > 0	•02	• 00	•	•	•00	• 70	1.00	1.00				
													
o l	• 47	.43	•39	.31	• 2 d	.29	-18	•12	-01	04	07	08	09
10	. 34	.27	.28	.20	.18	.20	•12	.06	02	06	09	10	10
20	03	04	03	04	03	•01	05	08	13	14	15	14	14
33	64	75	48	51	42	39	29	30	27	29	26	22	21
45	-Z.OH	-1.71	-1.49	-1.42	-1.32	96	82	74	59	50	43	38	32
50	-3.10	-2.08	-2.51	-2.25	-2.00	-1.74	-1.45	-1.16	88	75	62	57	47
75	-2.14	-3.40	-2.98	-2.76	-2.44	-2.25	-1.90	-1.65	-1.31	93	85	76	65
90	-+.20	-3.02	-3.48	- 3.05	-2.07	-2.40	-2.23	-1.79	-1-44	-1.19	-1.04	80	70
135	17.د-	-3.01	-2.98	-2.98	-2.52	-2.43	-1.77	-1.71	-1.40	-1.08	-1.00	91	78
120	-2.72	-2.18	-2.21	-2.00	-1.83	-1.87	-1.69	-1.30	-1.21	89	79	69	66
135	-2.10	-1.97	-1.81	-1.67	-1.62	-1.53	-1.30	-1.22	84	70	63	62	44
150	-1.70	-1.83	-1.80	-1.74	-1.49	-1.25	-1.23	-1.15	79	65	52	49	40
160	-1.40	-1.53	-1.62	-1.55	-1.47	-1.35	-1.17	86	67	52	40	35	31
170	-1.26	-1.31	-1.43	-1.47	-1.43	-1.24	-1.09	92	68	53	41	36	34
180	-1.21	-1.33	-1.38	-1.43	-1.35	-1.30	-1.16	91	65	52	41	35	29
-170	-l.30	-1.42	-1.45	-1.45	-1.30	-1.22	-1.08	95	66	49	40	35	25
-120		-2.23		-2.15		-1.88		-1.36		48		85	
-90		-3.97		-3.69		-2.86		-1.93		-1.28		85	
-60		-2.64		-2.00		-1.66		-1.21		73		56	
-30		30		23		10		19		20		19	
-10	.47	٠ ځو	• 30	.32	.29	.23	.17	.09	01	05	08	10	11
<u> </u>													
A r/D													
	د 2 · 2 غ	2.50	2 .7 5	3.30	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6.00	7.00
(deg)											<u> </u>		
1 .1				•									
0) 4	09	08	09	07	07	06	07	06	03	02		
10	10	Lu	09	08	07	08	06	08	06	04	03		
20	13	12	12	11	10	09	07	09	07	04	05		
اند	19	16	15	14	12	1.1	10	09	03	05	06	-01	
45	29	25	22	19	17	10	14	13	10	09	08	04	05
60	+4	30	32	28	20	22	20	16	12	11	10	06	
75	-•3u	47	43	36	30	28	20	20	14	16	13	09	
90													
1	~.63	50	49	43	38	32	31	25	23	19	16		
105	36	50 57	49 48	43 43	38 40	32 38	31 33	25 29	23 24	19 21	16 18	14	
120	36 54	50 57 50	49 48 44	43 43 3d	38 40 35	32 38 35	31 33 29	25 29 26	23 24 22	19 21 21	16 18 16	12	
120 135	54 41	50 57 50 3a	49 48 44 34	43 43 3d 26	38 40 35 28	32 38 35 21	31 33 29 21	25 29 26 10	23 24 22 14	19 21 21 13	16 18 16 13		09
120 135 150	36 54 41 20	50 57 50 30 26	49 48 44 34 19	43 43 36 26 16	38 40 35 28 13	32 38 35 21 10	31 33 29 21 10	25 29 26 10 08	23 24 22 14 04	19 21 21 13 05	16 18 16 13 05	12	09
120 135 150 160	36 54 41 20 24	50 57 50 30 26 17	49 48 44 34 19	43 43 36 26 16	38 40 35 28 13	32 38 35 21 10 07	31 33 29 21 10 08	25 29 26 10 08 05	23 24 22 14 04 06	19 21 21 13 05 05	16 18 16 13 05 04	12	09
120 135 150 160 170	36 54 41 20 24 25	50 57 50 30 26 17	49 48 44 34 19 22	43 43 3d 26 16 15	38 40 35 28 13 14	32 38 35 21 10 07	31 33 29 21 10 08	25 29 26 10 08 05	23 24 22 14 04 06	19 21 21 13 05 05	16 18 16 13 05 04	12	
120 135 150 160 170 180	36 54 41 20 24 25 29	50 57 50 30 26 17 23	49 48 44 34 19 22 18	43 43 36 26 16 15 16	38 40 35 28 13 14 15	32 38 35 21 10 07 11	31 33 29 21 10 08 08	25 29 26 10 08 05 07	23 24 22 14 06 07 08	19 21 21 13 05 05 01	16 18 16 13 05 04	12	09
120 135 150 160 170 180 -170	36 54 41 26 24 25 29	50 57 50 30 26 17	49 48 44 34 19 22 18 21	43 43 3d 26 16 15	38 40 35 28 13 14 15 10	32 38 35 21 10 07	31 33 29 21 10 08 08 13	25 29 26 10 08 05	23 24 22 14 04 06	19 21 21 13 05 05	16 18 16 13 05 04	12	
120 135 150 160 170 180 -170	36 54 41 26 24 25 29 25	50 57 50 30 26 17 23	49484434192218212040	43 43 36 26 16 15 16	38 40 35 28 13 14 15 10	32 38 35 21 10 07 11	31 33 29 21 10 08 08 13 11	25 29 26 10 08 05 07	23 24 22 14 06 07 08	19 21 21 13 05 05 01	16 18 16 13 05 04	12	
120 135 150 160 170 180 -170 -120 -90	36 54 41 26 25 25 25 29	50 57 50 30 26 17 23	49 48 44 34 19 22 18 21 20 40	43 43 36 26 16 15 16	38 40 35 28 13 14 15 10 12 41	32 38 35 21 10 07 11	31 33 29 21 10 08 08 13 11 31	25 29 26 10 08 05 07	23 24 22 14 06 07 08	19 21 21 13 05 05 01	16 18 16 13 05 04	12	
120 135 150 160 170 180 -170 -120 -90 -60	36 54 41 25 25 25 59 59	50 57 50 30 26 17 23	4948443419221821204932	43 43 36 26 16 15 16	3840352813141510124139	32 38 35 21 10 07 11	3133292110080813113131	25 29 26 10 08 05 07	23 24 22 14 06 07 08	19 21 21 13 05 05 01	16 18 16 13 05 04	12	
120 135 150 160 170 180 -170 -120 -60 -30	36 54 41 24 25 29 25 59 60 40	50 57 50 30 26 17 23 21	49484434192218212040493215	43 43 36 16 15 15 17	38 40 35 28 13 14 15 10 12 41 39 24 12	32 38 35 21 10 07 11 11	31332921100808131131322010	25 29 26 10 08 05 07 09	23 24 22 14 04 06 97 08	1921211305050107	16 18 16 13 05 04	12	
120 135 150 160 170 180 -170 -120 -90 -60	36 54 41 25 25 25 59 59	50 57 50 30 26 17 23	4948443419221821204932	43 43 36 26 16 15 16	3840352813141510124139	32 38 35 21 10 07 11	3133292110080813113131	25 29 26 10 08 05 07	23 24 22 14 06 07 08	19 21 21 13 05 05 01	16 18 16 13 05 04	12	

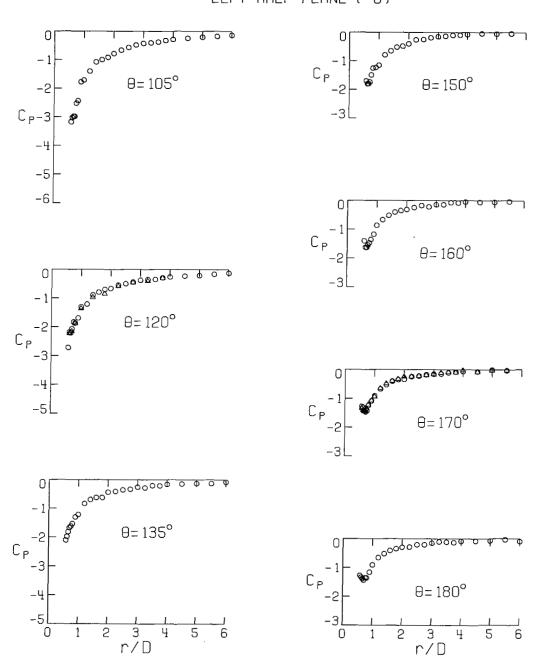
(a) C_p values.

Figure 9.- Pressure distribution on flat plate. R = 7.0; $M_j = 0.94$.



(b) Pressure distribution on rays.

Figure 9.- Continued.



(b) Concluded.

Figure 9.- Continued.

• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0) r/D=.58 C_P 0 -5 -6L -4 -5l C_P-2 -3 -4

(c) Pressure distribution on circles.

0

30

60

90

 Θ (DEGREES)

120

150 180

150

180

-5

0

30

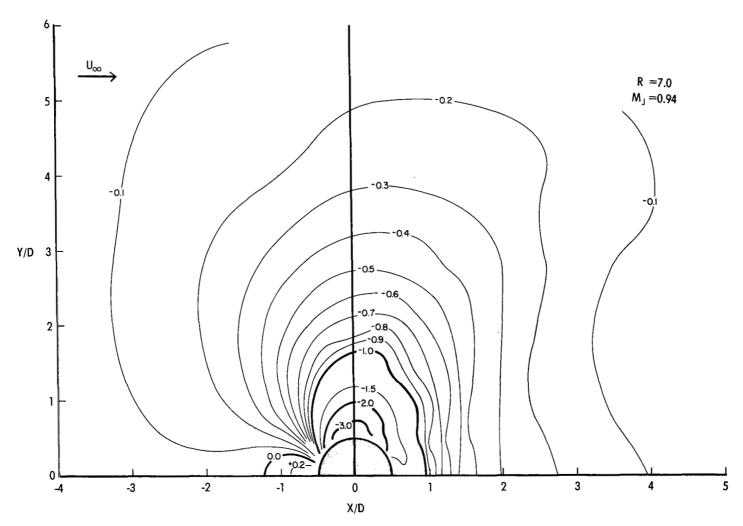
60

90

O(DEGREES)

120

Figure 9.- Continued.



(d) Contours of constant C_p .

Figure 9.- Concluded.

θ r/D (deg)	•5d	•02	• u6	.70	.75	.80	.90	1.00	1.20	1.40	1.60	1.80	2.00
	• 47	• 44	•40	٠.٤	.27	.22	.14	.06	04	08	12	14	13
10	•29	• 34	•24	.20	.14	.12	.07	.00	04	12	14	15	14
20	10	37	03	16	07	10	08	14	15	19	20	20	20
30	82	61	50	51	45	50	38	33	33	34	32	26	25
45	-2.24	-1.05	-1.42	-1.22	-1.35	98	80	76	54	58	45	42	38
60	-2.39	-2.43	-2.31	-2.13	-2.13	-1.63	~1.42	-1.10	95	74	72	62	53
75	-3.48	-2.93	-2.77	-2.46	-2.22	-2.22	-1.94	-1.41	-1.23	96	81	73	72
₹U	-3.85	-2.89	-2.66	-2.43	-2.43	-2.21	-1.88	-1.56	-1.31	-1.02	92	84	71
105	-2.34	-2.54	-2.30	-2.50	-2.02	-1.96	-1.81	-1.43	-1.11	-1.09	90	93	67
120	-1.71	-1.93	-1.84	-1.43	-1.80	-1.49	-1.70	-1.16	90	67	61	61	62
135	-1.59	-l.67	-1.60	-l.56	-1.24	-1.41	~1.05	-1.16	72	03	49	49	38
150	-1.47	-l.70	-1.61	-1.58	-1.2d	-1.09	-1.08	95	74	53	49	40	30
160	-1.14	-1.30	-1.43	-1.31	-1.19	-1.10	~1.02	78	54	42	31	35	26
176	-1.19	-1.2+	-1.18	-1.20	-1.14	-1.14	94	79	56	43	38	27	19
180	-1.37	-1.22	-1.19	-1.21	-1.17	-1.16	93	81	56	46	36	29	27
-170	-1.24	-1.11	-1.18	-1.29	-1.22	-1.21	94	76	55	37	32	30	30
-120		-2.05		-1.65		-1.54		-1.07		94		66	
-90		-3.41		-3.31		-2.76		-1.74		-1.18		90	
-60		-2.30		-2.30		-1.54		-1.05		72		56	
-30		35		23		22		26		22		25	
-10	• 46	• 39	-30	• 28	·24	-18	.11	.01	05	10	13	13	15
1 1													
											·		
θ r/D	2.25	2.20	2.75	3.00	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6.00	7.00
θ r/D (deg)	2.25	2.50	2.75	3.00	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6.00	7.00
10 /	2.25	2.50	2.75	3.00	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6.00	7.00
10 /	2.25	2.50	2.75	3.00	3.25	3.50	3.75	4.00	4.50 07	5.00	5.50	6.00	7.00
(deg) 0 10	13 14											6.00	7.00
(deg) 0 10 20	13 14 13	13 14 16	12 13 17	11 12 14	~.10 ~.12 ~.13	10 11 12	09 10 11	10 10 11	07 07 06	06 06 07	05 05 07	,	7.00
(deg) 0 10 20 30	13 14 16 21	-,13 -,14 -,16 -,21	12 13 17 18	11 12 14 18	~.10 ~.12 ~.13 ~.16	10 11 12 14	09 10 11 13	10 10 11 12	07 07 06 07	06 06 07 08	05 05 07 08	07	-
(deg) 0 10 20 30 45	13 14 16 21 34	13 14 16 21 28	12 13 17 18 28	11 12 14 18 24	10 12 13 16 21	10 11 12 14 18	09 10 11 13 17	10 10 11 12 15	07 07 06 07	06 06 07 08 11	05 05 07 08 10	07 09	7.00
(deg) 0 10 20 30 45 63	13 14 18 21 34 47	13 14 16 21 26 40	12 13 17 18 28 35	11 12 14 18 24 32	10 12 13 16 21 30	10 11 12 14 18 27	09 10 11 13 17	10 10 11 12 15 19	07 07 06 07 10	06 06 07 08 11 14	05 05 07 08 10	07 09 10	-
(deg) 0 10 20 30 45 60 75	13 14 18 21 34 47	13 14 16 21 26 40 49	12 13 17 18 28 35 43	11 12 14 18 24 32 40	10 12 13 16 21 30 35	10 11 12 14 18 27	09 10 11 13 17 23 28	10 10 11 12 15 19 20	07 07 06 07 10 11	06 06 07 08 11 14	05 05 07 08 10 12	07 09	-
(deg) 0 10 20 30 45 60 75	13 14 16 21 34 47 56	13 14 16 21 26 40 49	12 13 17 18 28 35 43 49	11 12 14 18 24 32 40	10 12 13 16 21 30 35 39	10 11 12 14 18 27 31	09 10 11 13 17 23 28 32	10 10 11 12 15 19 20 27	07 07 06 07 10 11 24	06 06 07 08 11 14 18	05 05 07 08 10 12 15	07 09 10 13	-
(deg) 0 10 20 30 45 60 75 90 105	13 14 13 21 34 47 56 53	13 14 10 21 28 40 49 54	12 13 17 18 28 35 43 49	11 12 14 18 24 32 40 48	~.10 ~.12 ~.13 ~.16 ~.21 ~.30 ~.35 ~.39 ~.42	10 11 12 14 18 27 31 37	0910111317232832	10 11 11 12 15 19 20 27	07 07 06 07 10 11 24 25	06 06 07 08 11 14 18 22	05 05 07 08 10 12 15 17	07 09 10 13	-
(deg) 0 10 20 30 45 60 75 90 105 120	13 14 13 21 34 47 56 59	13 14 16 21 28 49 49 54	12 13 17 18 28 35 43 49 43	11 12 14 18 24 32 40 48 41	10 12 13 16 21 30 35 39 42	10 11 12 14 18 27 31 37	0910111317232832	10 10 11 12 15 19 20 27 27	07 07 06 07 10 11 24 25 25	06 06 07 08 11 14 18 22 22	05 05 07 08 10 12 15 17 20	07 09 10 13 16	01
(deg) 0 10 20 30 45 62 75 90 105 120 135	13 14 13 21 34 47 50 58 59	13 14 16 21 28 40 49 54 54	12 13 17 18 28 35 43 49 43	11 12 14 18 24 32 40 48 41	10 12 13 16 21 30 35 39 42 34	10 11 12 14 18 27 31 37 37	091011131723283231	10 10 11 12 15 19 20 27 27	07 07 06 07 10 11 24 25 25 25	06 06 07 08 11 14 18 22 22 21	05 05 07 08 10 12 15 17 20 17	07 09 10 13	-
(deg) 0 10 20 30 45 60 75 90 105 120 135 150	13 14 13 21 34 47 56 53 59 46 30 20	13 14 16 21 28 40 49 54 54 42 27	12 13 17 18 28 35 43 49 43 36 30	1112141824324048413024	~.10 ~.12 ~.13 ~.16 ~.21 ~.35 ~.39 ~.42 ~.34 ~.25	10 11 12 14 18 27 31 37 39 32 22	09101113172328323128	10 13 11 12 15 19 20 27 29 22 15	07 07 06 07 10 11 24 25 25 22 16	06 06 07 08 11 14 18 22 22 21 12	05 05 07 08 10 12 15 17 20 17	07 09 10 13 16	01
(deg) 0 10 20 30 45 65 75 90 105 120 135 150 160	13 14 13 21 34 47 56 59 46 36 36 20 24	13 14 10 21 28 40 49 54 54 27 27	12 13 17 18 28 35 43 49 43 30 15	11 12 14 18 24 32 40 41 30 24 17	~.10 ~.12 ~.13 ~.16 ~.21 ~.30 ~.35 ~.39 ~.42 ~.34 ~.25 ~.13 ~.08	10 11 12 14 18 27 31 37 39 32 22 14	0910111317232831281707	10 11 11 12 15 19 20 27 29 22 15 07	07 07 06 07 10 11 24 25 25 25 26 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06	07 09 10 13 16	01
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170	13 14 13 21 34 47 56 53 59 46 36 26 24 25	13 14 16 21 28 40 49 54 54 27 20 17	12 13 17 18 28 35 43 49 43 36 30 15 19	111214182432404130241714	10 12 13 16 21 30 35 39 42 34 25 13	10 11 12 14 18 27 31 37 39 32 22 14	091011131723283231281707	10 11 12 15 19 20 27 29 22 15 07	07 07 06 07 10 11 24 25 25 22 16 04 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06 01	07 09 10 13 16 16	01
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180	13 14 13 21 34 47 56 59 46 30 26 24 25 20	131416212849495454272017	12 13 17 18 28 35 43 43 36 30 15 19 17	1112141824324041302417114	10 12 13 16 21 30 35 39 42 34 25 13 08	10 11 12 14 18 27 31 37 39 32 22 14 08 09	091011131723283128170708	10 11 12 15 19 20 27 29 22 15 07 07	07 07 06 07 10 11 24 25 25 22 16 04 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06	07 09 10 13 16	01
(deg) 0 10 20 30 45 62 75 90 105 120 135 150 160 170 180 -170	13 14 13 21 34 47 56 59 46 30 26 25 25 25	13 14 16 21 28 40 49 54 54 27 20 17	12 13 17 18 28 35 43 49 43 36 36 31 15 19	111214182432404130241714	~.10 ~.12 ~.13 ~.16 ~.21 ~.30 ~.35 ~.39 ~.42 ~.34 ~.25 ~.13 ~.08	10 11 12 14 18 27 31 37 39 32 22 14	0910111317232832312817070808	10 11 12 15 19 20 27 29 22 15 07	07 07 06 07 10 11 24 25 25 22 16 04 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06 01	07 09 10 13 16 16	01
(deg) 0 10 20 30 45 65 75 90 105 120 135 150 160 170 180 -170 -123	13 14 13 21 34 47 56 59 46 30 26 24 25 25 25	131416212849495454272017	12 13 17 18 28 35 43 49 43 36 30 15 19 17	1112141824324041302417114	10121316213035394234251308111608	10 11 12 14 18 27 31 37 39 32 22 14 08 09	09101113172328323128070708080810	10 11 12 15 19 20 27 29 22 15 07 07	07 07 06 07 10 11 24 25 25 22 16 04 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06 01	07 09 10 13 16 16	01
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120 -90	13 14 13 21 34 47 56 59 46 36 29 24 25 25 24	131416212849495454272017	12 13 17 18 28 35 49 43 30 15 19 17 18 17	1112141824324041302417114	~.10 ~.12 ~.13 ~.16 ~.21 ~.30 ~.35 ~.39 ~.42 ~.34 ~.25 ~.13 ~.08 ~.11 ~.08 ~.08	10 11 12 14 18 27 31 37 39 32 22 14 08 09	0910111317233128170708081024	10 11 12 15 19 20 27 29 22 15 07 07	07 07 06 07 10 11 24 25 25 22 16 04 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06 01	07 09 10 13 16 16	01
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120 -90 -60	13141321344756594635202425202347	131416212849495454272017	12 13 17 18 28 35 43 43 36 30 15 17 18 17	1112141824324041302417114	1012131621303539423425111608414143	10 11 12 14 18 27 31 37 39 32 22 14 08 09	0910111317232832312817070808102434	10 11 12 15 19 20 27 29 22 15 07 07	07 07 06 07 10 11 24 25 25 22 16 04 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06 01	07 09 10 13 16 16	01
(deg) 0 10 20 30 45 60 75 90 105 120 135 150 160 170 180 -170 -120 -90	13 14 13 21 34 47 56 59 46 36 29 24 25 25 24	131416212849495454272017	12 13 17 18 28 35 49 43 30 15 19 17 18 17	1112141824324041302417114	~.10 ~.12 ~.13 ~.16 ~.21 ~.30 ~.35 ~.39 ~.42 ~.34 ~.25 ~.13 ~.08 ~.11 ~.08 ~.08	10 11 12 14 18 27 31 37 39 32 22 14 08 09	0910111317233128170708081024	10 11 12 15 19 20 27 29 22 15 07 07	07 07 06 07 10 11 24 25 25 22 16 04 04	06 06 07 08 11 14 18 22 22 21 12 04 03	05 05 07 08 10 12 15 17 20 17 11 06 01	07 09 10 13 16 16	01

(a) C_p values.

Figure 10.- Pressure distribution on flat plate. R = 8.0; $M_j = 0.93$.

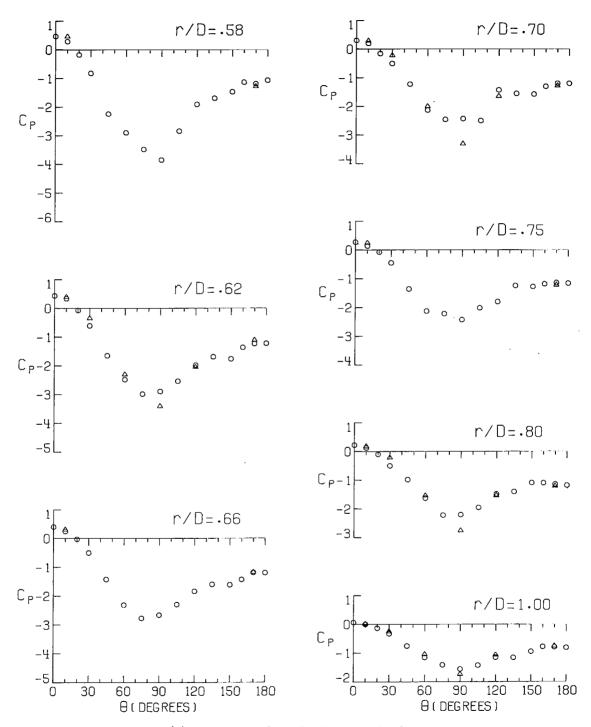
• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0) C_P 0 0=0° 0=10° 0=20°

(b) Pressure distribution on rays.

Figure 10.- Continued.

• RIGHT HALF-PLANE (0) • LEFT HALF-PLANE (-0) -6 -5[r/D (b) Concluded.

Figure 10.- Continued.



(c) Pressure distribution on circles.

Figure 10.- Continued.

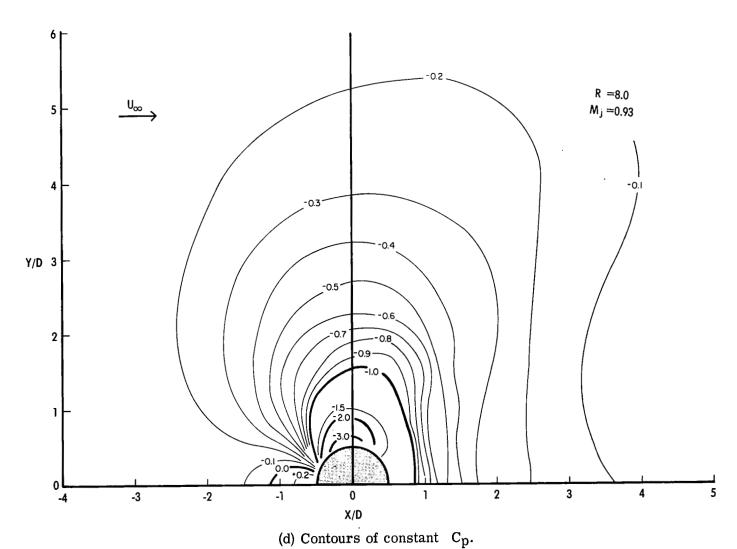


Figure 10.- Concluded.

θ r/D (deg)	•5₫	•62	.66	.70	.75	.80	•90	1.00	1.20	1.40	1.60	1.80	2.00
3	• 41	•4J	. 34	•26	.21	.19	.09	.01	09	14	17	18	17
10	•21	.22	. 25	-10	.08	•06	01	07	10	18	19	19	17
20	19	.01	13	i s	17	10	16	20	23	22	23	21	24
30	70	53	54	35	39	39	32	37	31	35	33	28	29
45	-1.69	-1.3+	-1.33	-1.34	-1.16	90	73	75	61	45	47	40	42
60	-2.49	-2.35	-2.06	-1.92	-1.40	-1.47	-1.10	-1.23	82	75	63	52	54
75	13 و د	-2.74	-2.74	-2.09	-2.11	-1.68	-1.40	-1.25	-1.00	82	61	73	59
9.1	- 2.7੪	-2.80	-2.28	-2.32	-1.95	-1.67	-1.75	-1.36	94	89	79	73	57
105	-2.10	~2.03	-2.03	-1.54	-1.53	-1.56	-1.35	-1.06	95	61	73	72	60
120	-1.51	-1.37	-1.22	-1.43	47	-1.50	-1.31	91	74	75	46	52	38
135	-1.50	-1.27	-1.25	-1.16	-1.30	-1.09	-1.06	49	68	55	40	30	25
150	-1.09	96	-1.28	-1.16	-1.05	71	76	70	50	42	32	28	23
160	ده	-1.17	98	85	-1.01	84	06	54	36	32	22	22	19
170	37	97	-1.01	85	93	83	73	56	34	31	21	19	17
180	84	93	94	96	-1.00	91	74	54	40	33	23	18	13
-170	-1.31	yó	98	-1.33	47	89	82	04	39	30	28	16	17
-120		-1.75		-1.63		-1.70		-1.03		77		70	
- 90		-3.20		-2.70		-2.19		-1.43		97		69	
-60		-2.80		-1.99		-1.61		-1.02		59		56	
-30		34		35		30	_	34		32		28	
-10	• 34	-3L	•27	• 25	.19	.13	•06	•00	10	15	18	17	21
										,-			
θ r/D (deg)	2.25	2.50	2.75	3.00	3.25	3.50	3.75	4.00	4.50	5.00	5.50	6.00	7.00
(deg)						<u> </u>						6.00	7.00
(deg)	1+	19	18	16	14	14	14	32	09	10	10	6.00	7.00
(deg)	1 # 21	19 20	18 18	16 18	14 17	14 14	14 15	32 31	09 10	10 11	10 09	6.00	7.00
(deg) 0 10 20	1+ 21 23	19 20 23	18 16 20	16 18 20	14 17 17	14 14 16	14 15 15	32 31 29	09 1J 11	10 11 11	10 09 10		7.00
(deg) 0 10 20 30	1 + 21 23 28	19 20 23 24	18 16 20 22	16 18 20 23	14 17 17 20	14 14 16 20	14 15 15 19	32 31 29	09 10 11 12	10 11 11	10 09 10 11	08	
(deg) 10 20 30 40	1 + 21 23 28	19 20 23 24 30	18 16 20 22 30	16 18 20 23 29	14 17 17 20 27	14 14 16 20 24	14 15 15 19 23	32 31 29 29 28	09 10 11 12 14	10 11 11 12 16	10 09 10 11 12	08 09	7.00
(deg) 0 10 20 30 40 60	1 / 21 23 28 30 40	19 20 23 24 30 43	18 16 20 22 30 31	16 18 20 23 29 32	14 17 17 20 27 31	14 14 16 20 24	14 15 15 19 23 27	32 31 29 29 28 28	09 10 11 12 14 19	10 11 11 12 16 17	10 09 10 11 12 15	08 09 14	
(deg) 10 20 30 40	1 / 21 23 28 30 40	19 20 23 24 30 43	18 16 20 22 30 31	16 18 20 23 29 32	14 17 17 20 27 31	14 14 16 20 24 26	14 15 15 19 23 27	32 31 29 29 28 28	09 10 11 12 14 19	10 11 11 12 16 17	10 09 10 11 12 15	08 09	
(deg) 10 20 30 40 60 75	1 / 21 23 28 30 40 40	19 20 23 24 30 43 50	18 16 20 22 30 31 42 45	16 18 20 23 29 32 38 34	14 17 17 20 27 31 33 39	14 14 16 20 24 26 31	14 15 15 19 23 27 30	32 31 29 29 28 28 29	09 10 11 12 14 19 21	10 11 11 12 16 17 21	10 09 10 11 12 15 18	08 09 14 16	
(deg) 10 20 30 40 60 75	1 / 21 23 28 30 40 54	19 20 23 24 30 43	18 16 20 22 30 31	16 18 20 23 29 32	14 17 17 20 27 31 33 39	14 14 16 20 24 26 31 40	14 15 15 19 23 27 30 35	32 31 29 29 28 28	09 10 11 12 14 19	10 11 11 12 16 17	10 09 10 11 12 15	08 09 14	
(deg) 10 20 30 40 60 75 90 105	1 / 21 23 28 30 40 54 54	19 20 23 24 30 43 50 51	18 18 20 22 30 31 42 45 38	16 18 20 23 29 32 32 34	14 17 17 20 27 31 33 39	14 14 16 20 24 26 31	14 15 15 19 23 27 30	32 31 29 29 28 24 29 29	09 10 11 12 14 19 21	10 11 11 12 16 17 21 21	10 09 10 11 12 15 18 20	08 09 14 16	
(deg) 0 10 20 30 40 60 75 90 105 120	1 / 21 23 28 30 40 54	19 20 23 24 30 43 50 49	18 18 20 22 30 31 42 45 38	16 18 20 23 29 32 38 34 41	14 17 17 20 27 31 33 39 32	14 14 16 20 24 26 31 40 29	14 15 15 19 23 27 30 35 27	32 31 29 29 28 28 29 30 29	091011121419212121	10 11 11 12 16 17 21 21	10 09 10 11 12 15 18 20 17	08 09 14 16	01
(deg) 10 20 30 40 60 75 90 105 120 130	1 y212328304040545454	19 20 23 24 30 43 50 49 51	18 16 20 22 30 31 42 45 38 33 16	16 18 20 23 32 32 38 34 41 29	14 17 17 20 27 31 33 39 32 25	14 16 20 24 26 31 40 29 28	14 15 15 19 23 27 30 35 27	32 31 29 29 28 29 29 30 29	09131112141921212121	10111112161721212010	10 09 11 12 15 18 20 17 14	08 09 14 16	01
(deg) 10 20 30 40 60 75 90 105 120 130 150	17 21 23 28 30 40 54 54 54 25	19 20 23 24 30 43 50 49 51 30	18 20 22 30 31 42 45 38 33	16 18 20 23 32 32 34 41 29 09	14 17 17 20 27 31 33 39 32 25 16	14 16 20 24 26 31 40 29 28 15	14 15 15 19 23 27 30 35 27 24 12	32 31 29 29 28 26 29 30 29	0910111214192121212100	1011111216172121201609	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16	01
(deg) 0 10 20 30 40 60 75 90 105 120 130 150 160	17 21 23 28 30 40 54 54 55 20 20	19202324304351302016	18162022303142453833161713	1618202329323441291919	14 17 17 20 27 31 33 39 32 25 16 10	14 16 20 24 26 31 40 29 28 15	14 15 15 19 23 27 30 35 27 24 12	32 31 29 29 28 29 29 30 29 10	091J1112141921212121060402	101112161721212010090301	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16	01
(deg) 0 10 20 30 40 60 75 90 105 120 130 150 150 160 170	1 / 21 23 30 40 54 54 54 20 20 10	19202324304351302615	1818202230314245383316171311	16 18 20 23 29 32 34 41 29 19 09	14 17 17 20 27 31 33 39 32 25 16 10	14 16 20 24 26 31 40 29 28 15 14	14 15 15 19 23 27 30 35 27 24 12 07	32 31 29 28 28 29 30 29 10	09101112141921212121000402	101112161721212010090301	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16 18 11	01
(deg) 0 10 20 30 40 60 75 90 105 120 130 150 160 170 180	1 y21232830404054545220181913	1920232430435130261615	1818202230314245383316171311	16 18 20 23 32 32 34 41 29 19 09	14 17 17 20 27 31 33 39 32 25 16 10 10	14 16 20 24 26 31 40 29 28 15 14 12	14 15 15 19 23 27 30 35 27 24 12 07 10	32 31 29 28 28 29 30 29 10 08 10	091011121419212121210004	10111216172120160903010406	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16 18 11	01
(deg) 10 20 30 40 60 75 90 105 120 130 150 160 170 180 -170	1 y2123283040405454542018191314	1920232430435130261615	182022303142453833161713111208	16 18 20 23 32 32 34 41 29 19 09	14 17 17 20 27 31 33 39 32 25 16 10 10	14 16 20 24 26 31 40 29 28 15 14 12	14 15 15 19 23 27 30 35 27 24 12 07 10	32 31 29 28 28 29 30 29 10 08 10	091011121419212121210004	10111216172120160903010406	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16 18 11	01
(deg) 10 20 30 40 60 75 90 105 120 130 150 160 170 180 -170 -120	17212328304054545920181913	1920232430435130261615	18202230314245383316171311120836	16 18 20 23 32 32 34 34 41 29 19 09	14 17 17 20 27 31 33 39 32 25 16 10 10 04 13	14 16 20 24 26 31 40 29 28 15 14 12	14 15 19 23 27 30 35 27 12 07 10 03 14	32 31 29 28 28 29 30 29 10 08 10	091011121419212121210004	1011121617212120160903010406	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16 18 11	01
(deg) 0 10 20 30 40 60 75 90 105 120 130 150 160 170 180 -170 -120 -90	17212328304054552018191314	1920232430435130261615	1820223031424538331617131112083645	16 18 20 23 32 32 34 34 41 29 19 09	1417202731333932251610041307	14 16 20 24 26 31 40 29 28 15 14 12	14 15 19 23 27 30 27 24 12 07 10 03 14 03 29	32 31 29 28 28 29 30 29 10 08 10	091011121419212121210004	1011121617212120160903010406	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16 18 11	01
(deg) 0 10 20 30 40 60 75 90 105 120 130 150 170 180 -170 -120 -90 -60	17 21 23 30 40 54 54 59 20 10 19 13 35 35 47	1920232430435130261615	1818202230314238331617131112083639	16 18 20 23 32 32 34 34 41 29 19 09	14171720273133393225161010041307343935	14 16 20 24 26 31 40 29 28 15 14 12	14 15 19 23 27 30 35 27 24 12 07 10 03 14 03	32 31 29 28 28 29 30 29 10 08 10	091J1112141921212104040402	1011121617212120160903010406	10 09 10 11 12 15 18 20 17 14 07 00	08 09 14 16 18 11	01

(a) C_p values.

Figure 11.- Pressure distribution on flat plate. R = 10.0; $M_j = 0.93$.

C_P 0 0=0° 0=10° C_P 0

(b) Pressure distribution on rays.

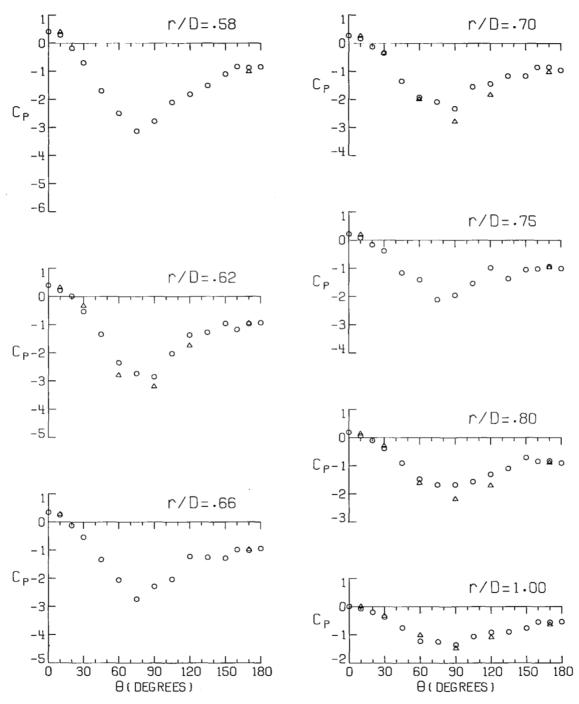
Figure 11.- Continued.

(b) Concluded.

Figure 11.- Continued.

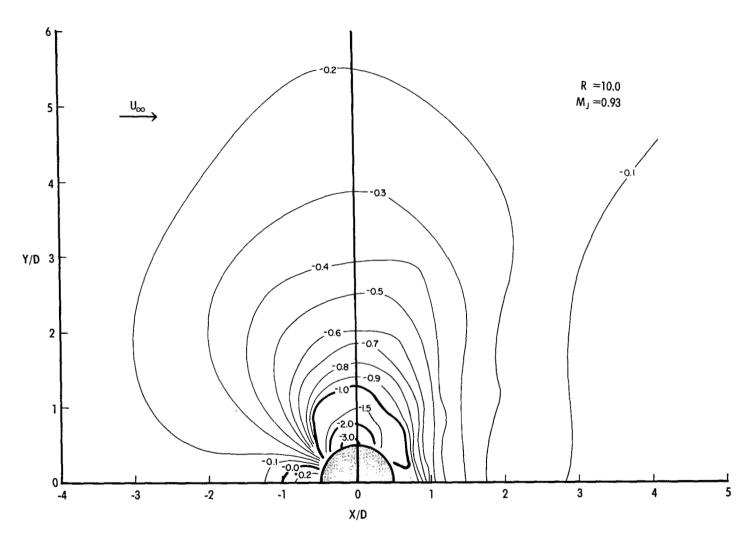
-5

3 r/D



(c) Pressure distribution on circles.

Figure 11.- Continued.



(d) Contours of constant C_p .

Figure 11.- Concluded.

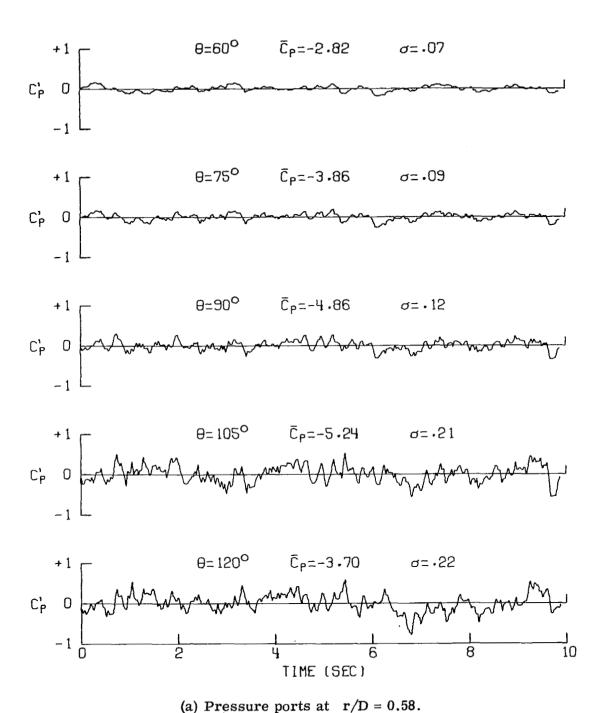


Figure 12.- Time variation of C_p with pressure-transducer output filtered. $R=4.1;\ M_{\dot 1}=0.48.$

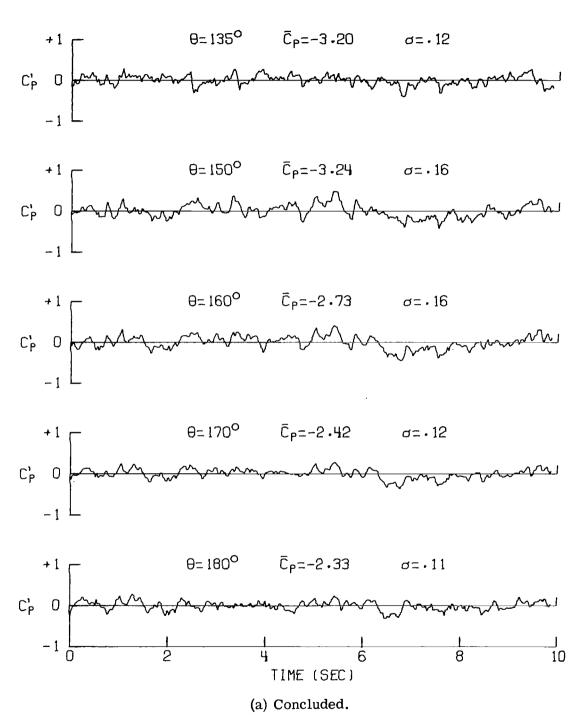


Figure 12.- Continued.

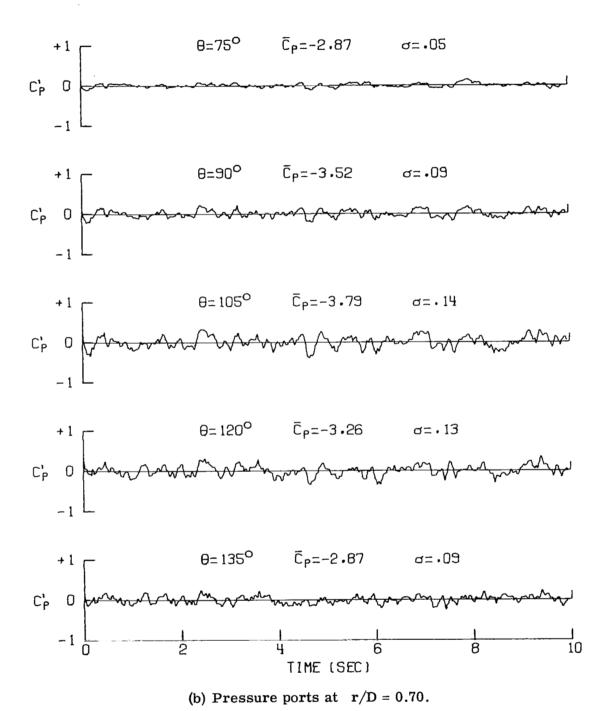


Figure 12.- Continued.

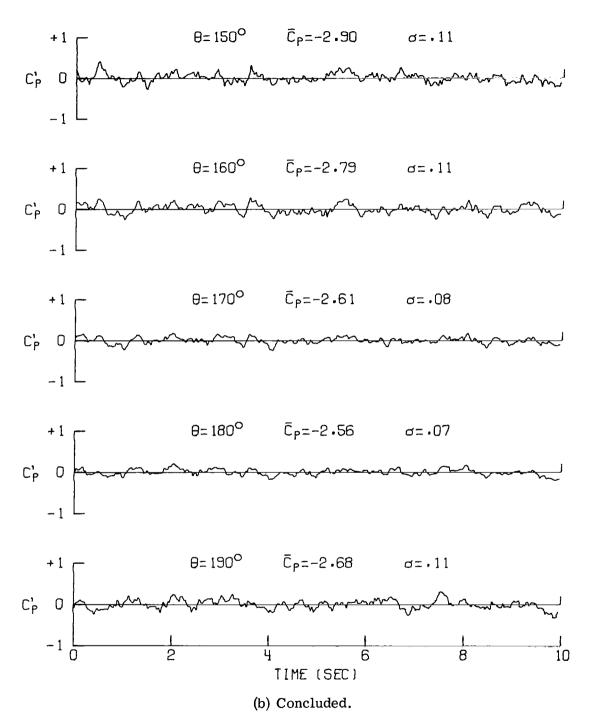


Figure 12.- Continued.

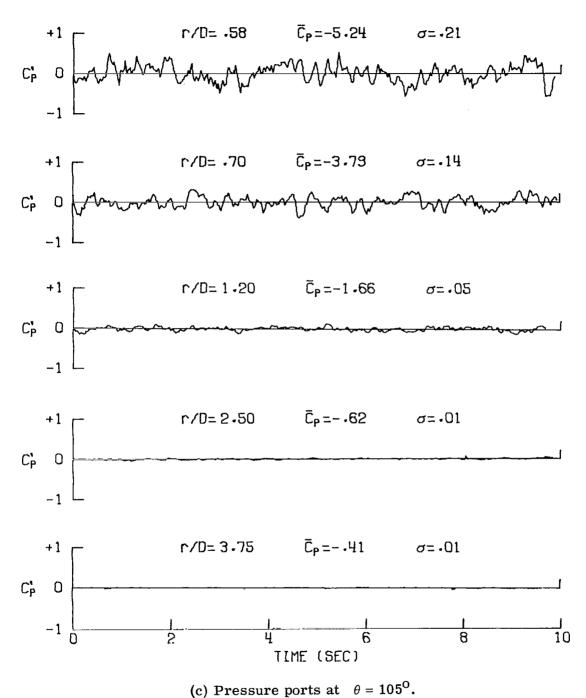
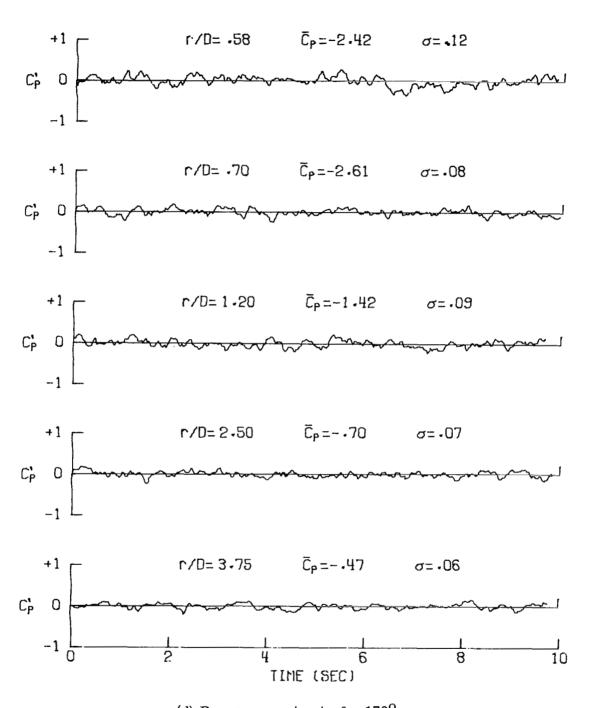


Figure 12.- Continued.



(d) Pressure ports at $\theta = 170^{\circ}$. Figure 12.- Concluded.

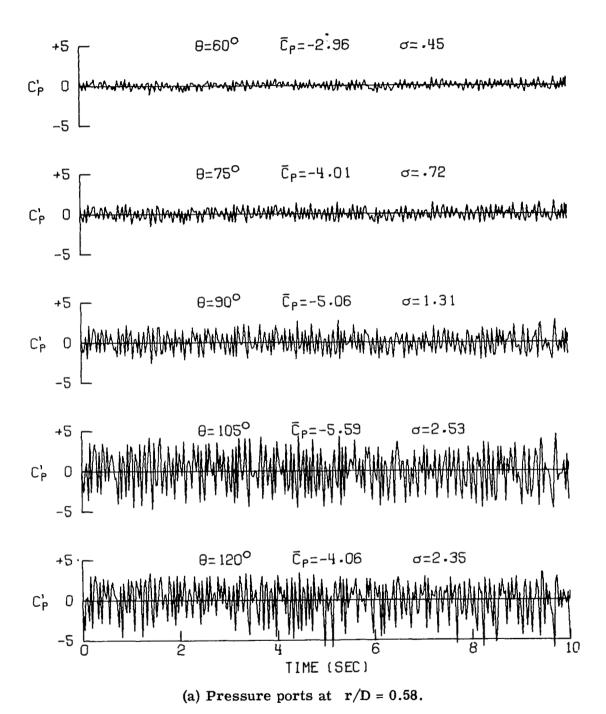


Figure 13.- Time variation of C_p with pressure-transducer output not filtered. $R=4.1;\ M_j=0.48.$

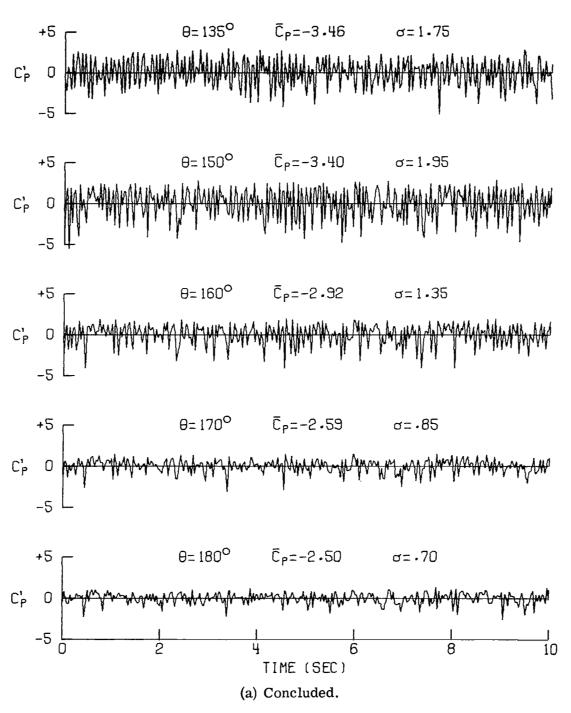
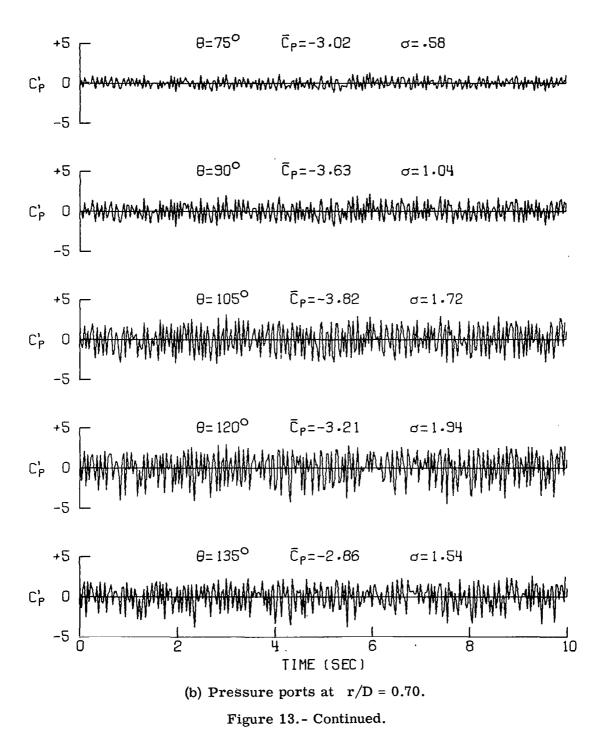


Figure 13.- Continued.



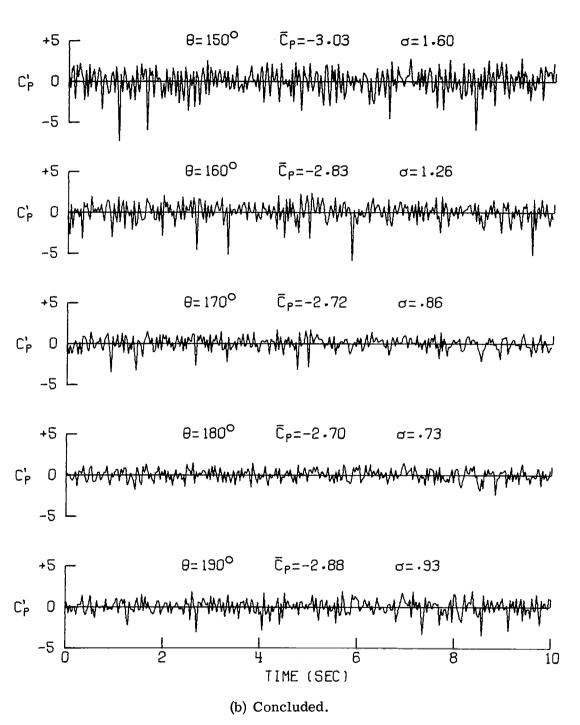
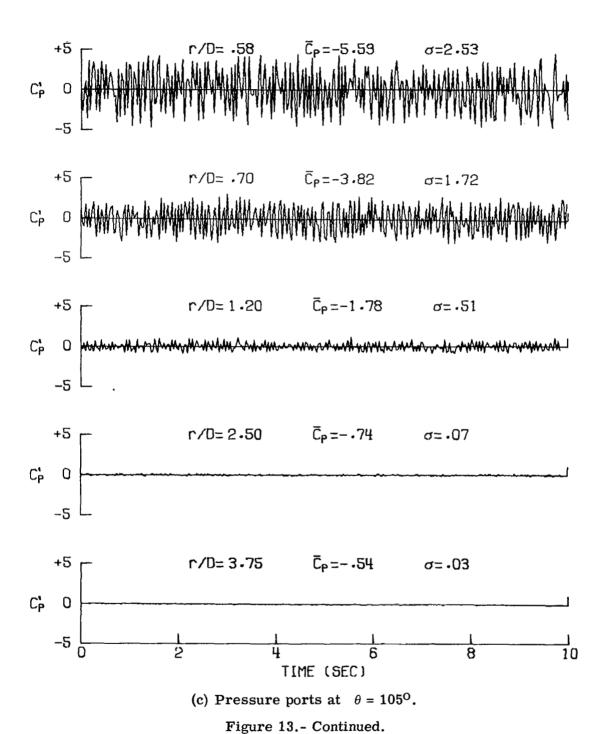
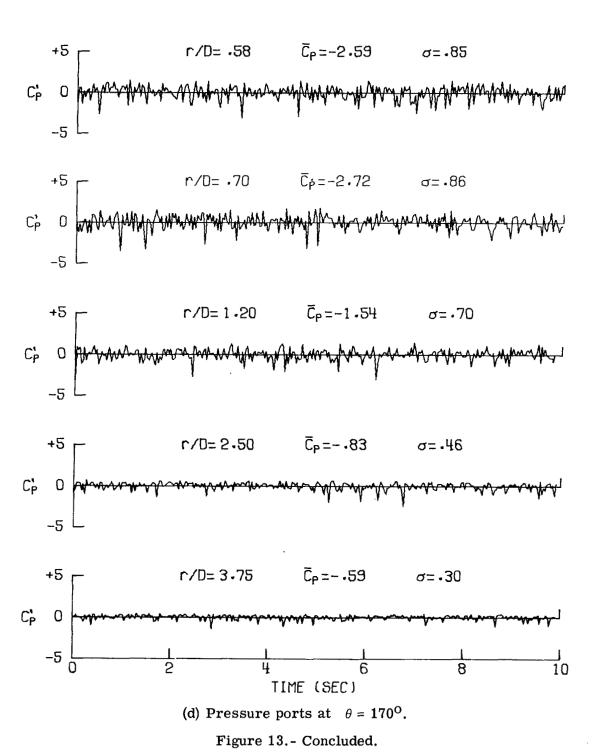


Figure 13.- Continued.





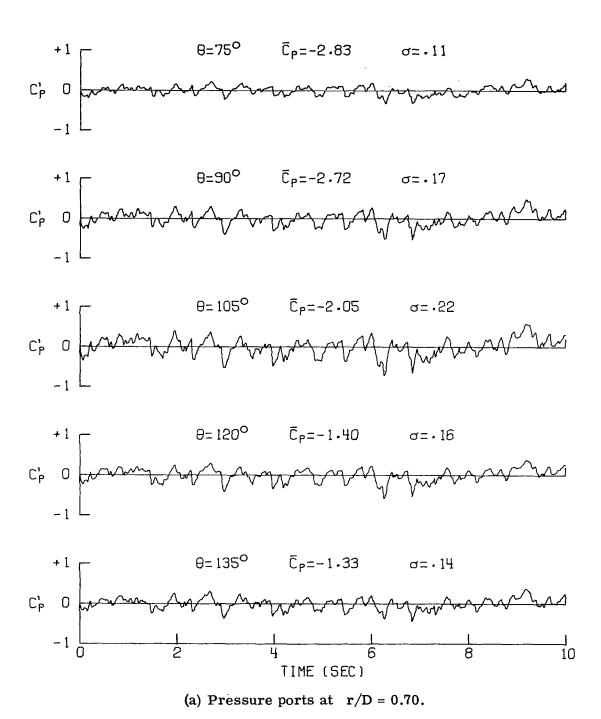


Figure 14.- Time variation of C_p with pressure-transducer output filtered. R = 8.0; $M_{\mbox{\scriptsize j}}$ = 0.93.

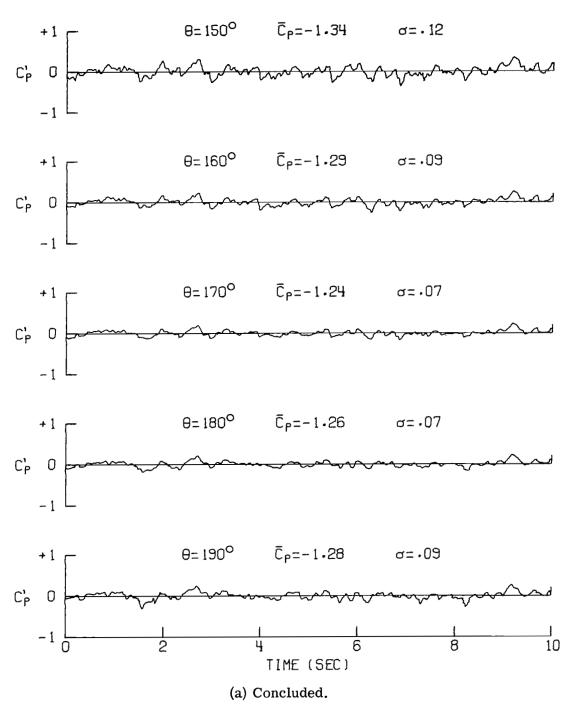
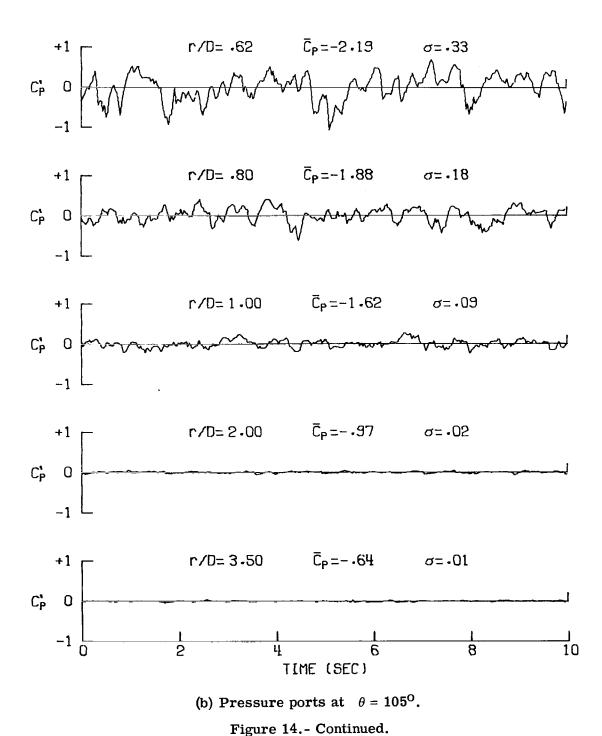
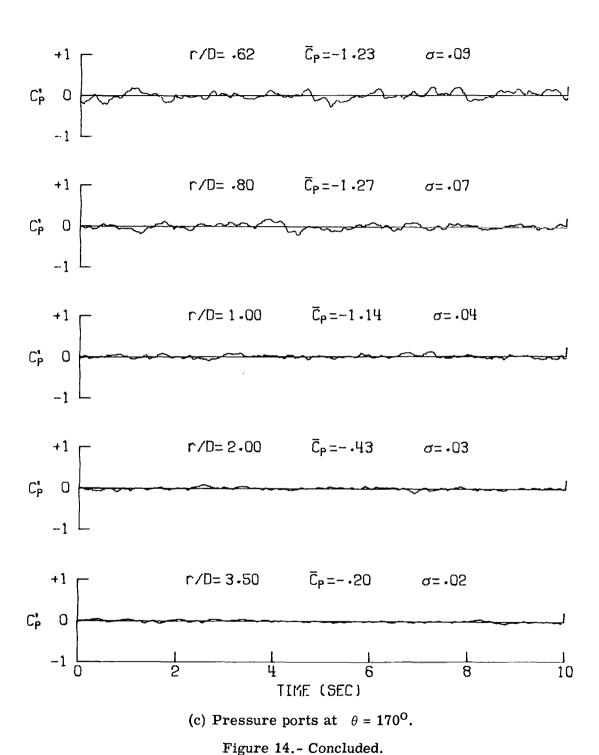


Figure 14.- Continued.





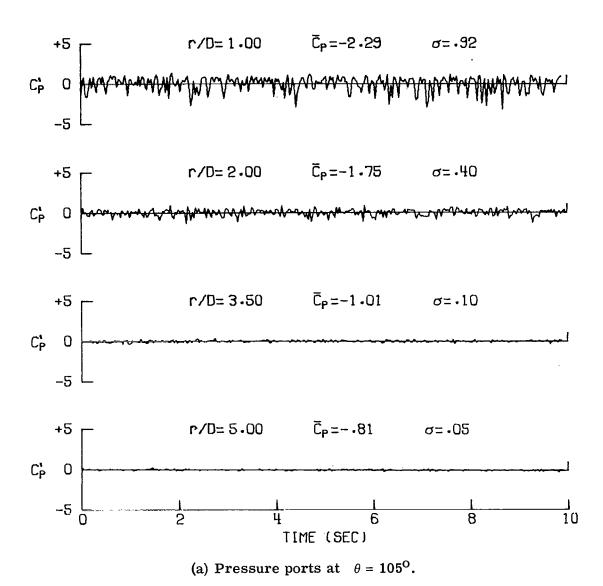


Figure 15.- Time variation of $\, C_p \,$ with pressure-transducer output not filtered. $\, R = 8.0; \,$ $\, M_j = 0.93.$

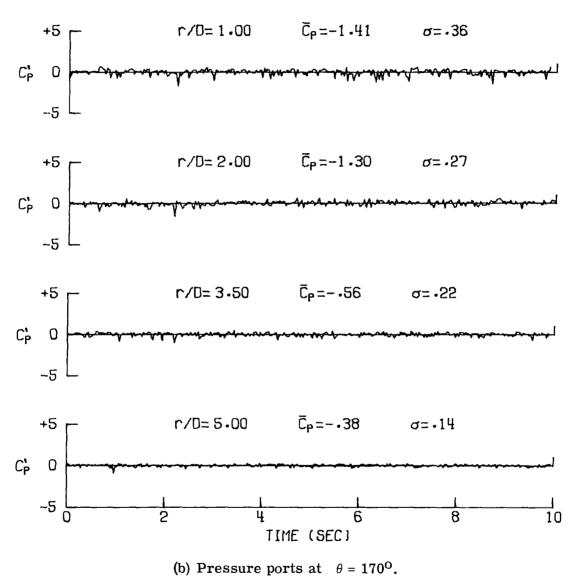


Figure 15.- Concluded.

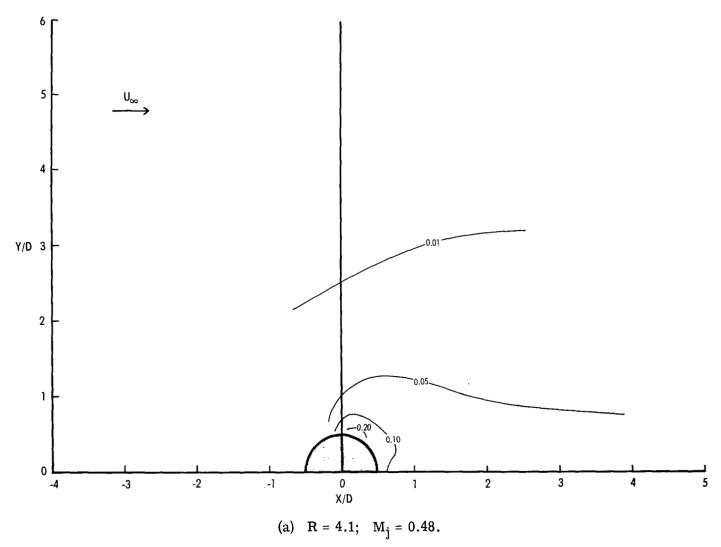
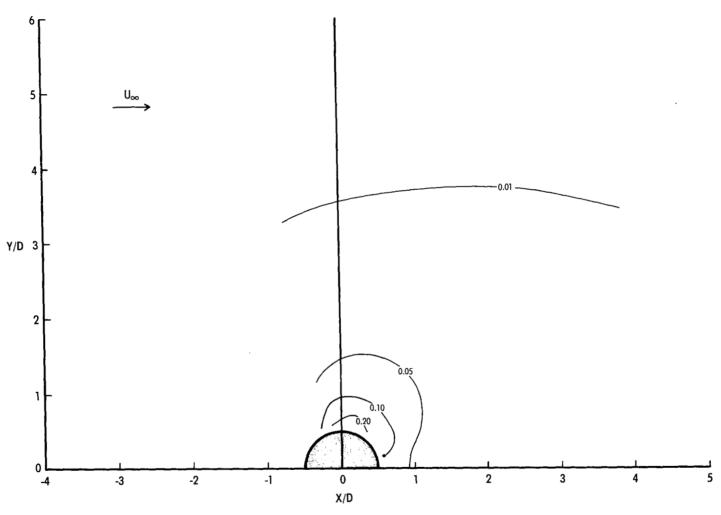
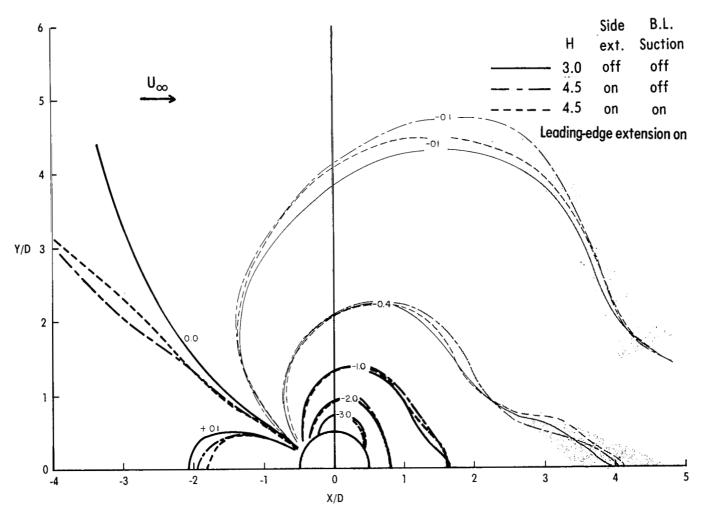


Figure 16.- Contours of constant σ .



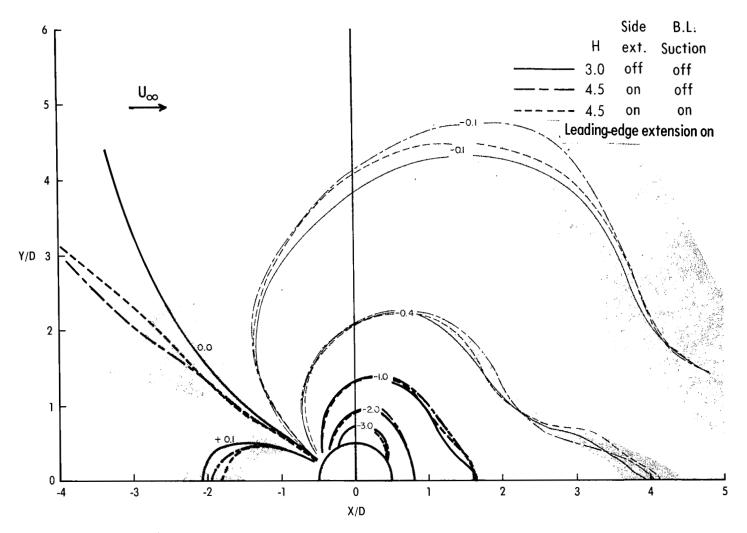
(b) R = 8.0; $M_j = 0.93$.

Figure 16.- Concluded.



(a) Uncertainty due to measured fluctuations in $\ C_{p}.\ \Delta C_{p}$ = $\sigma.$

Figure 17.- Effect of test conditions (see table 1) on pressure distribution and uncertainty bands.



(b) Uncertainty due to pressure-transducer errors. ΔC_p = 0.036.

Figure 17.- Concluded.

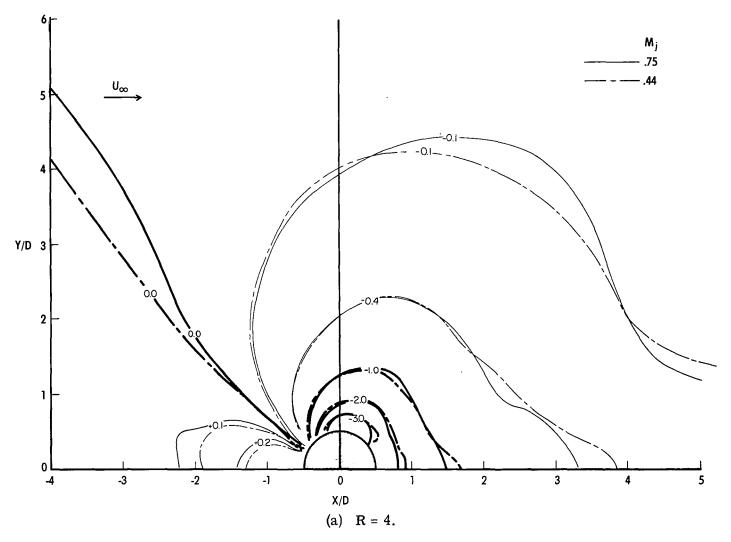


Figure 18.- Effect of changing jet Mach number on pressure distribution.

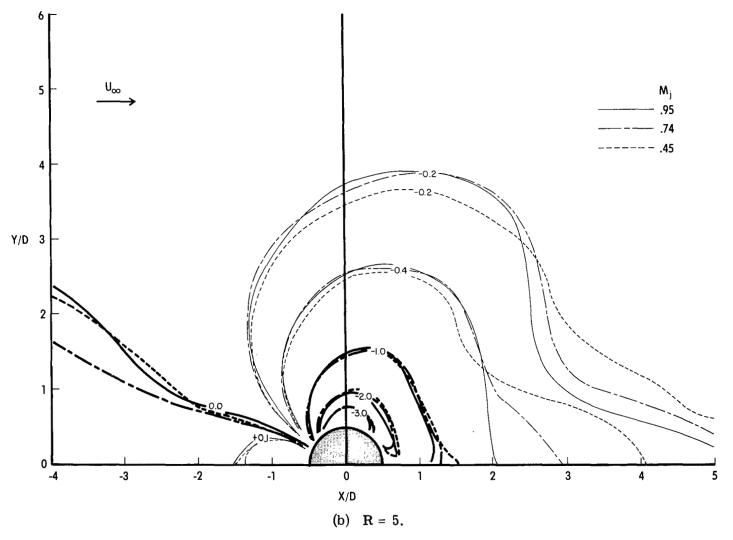


Figure 18.- Continued.

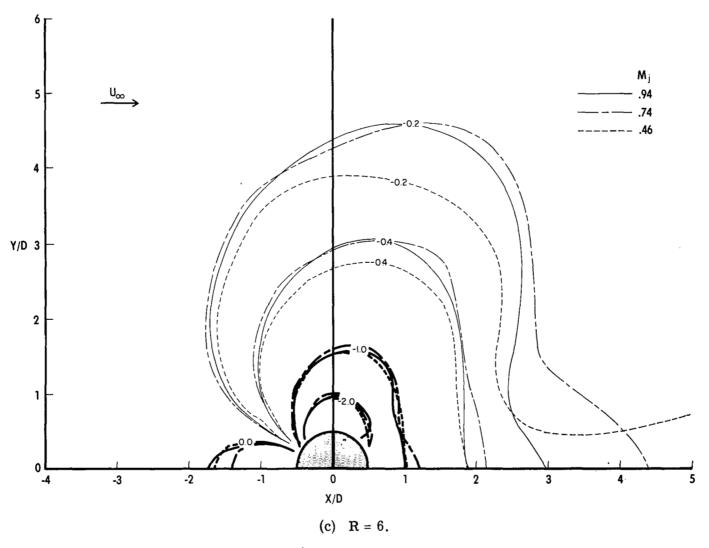


Figure 18.- Continued.

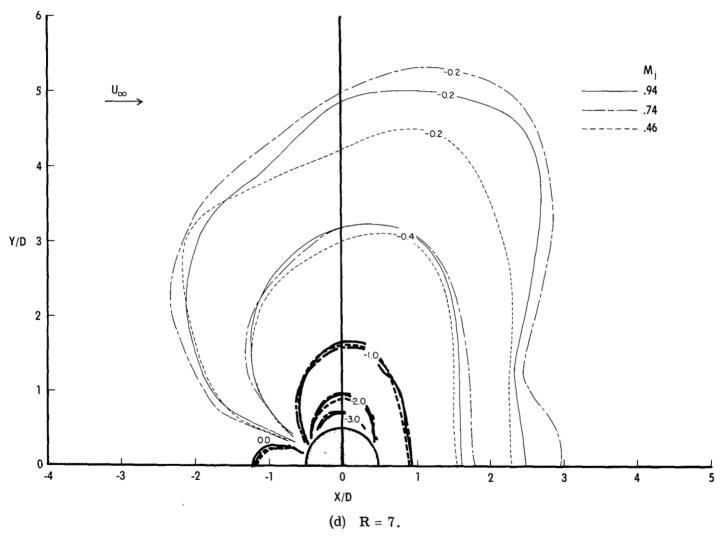


Figure 18.- Continued.

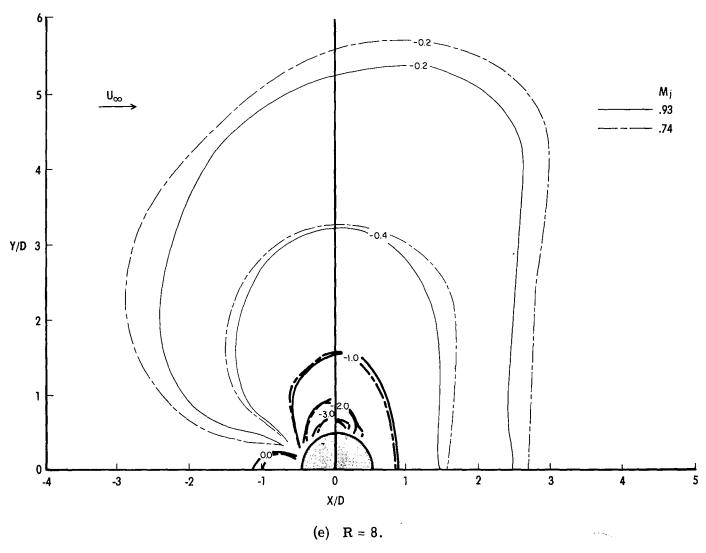


Figure 18.- Continued.

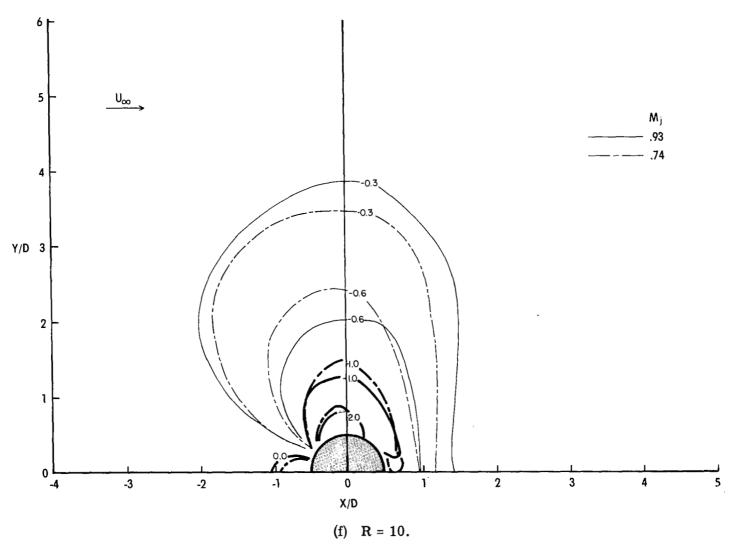


Figure 18.- Concluded.

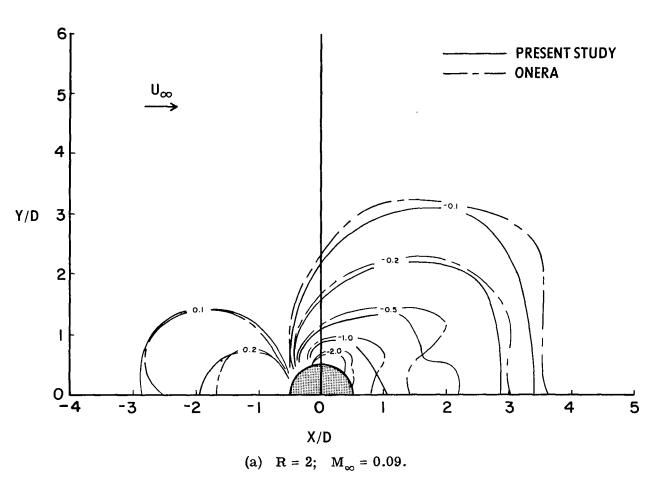


Figure 19.- Comparison with Soullier (ref. 21).

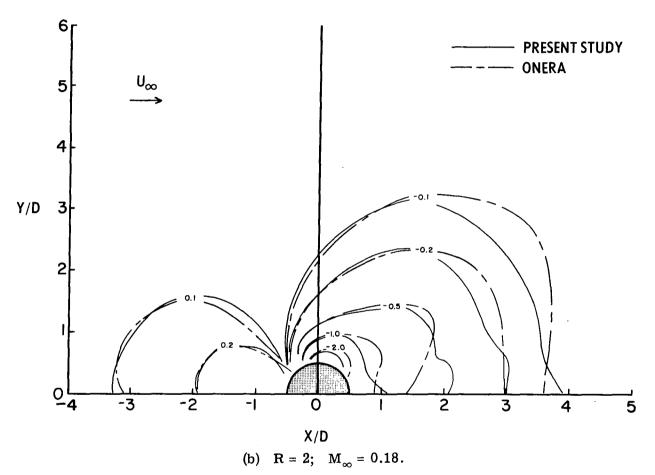


Figure 19.- Continued.

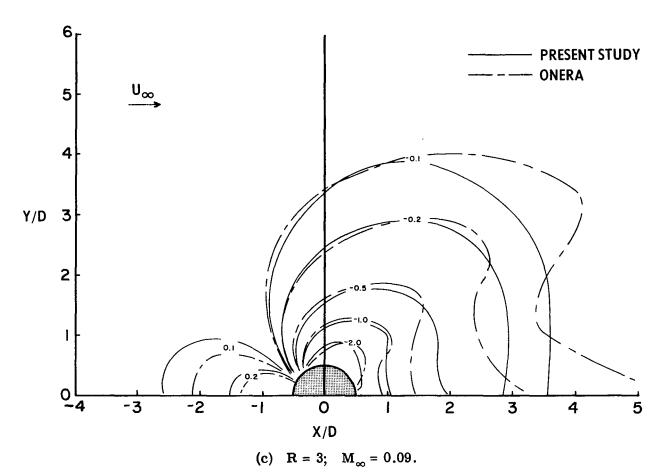


Figure 19.- Continued.

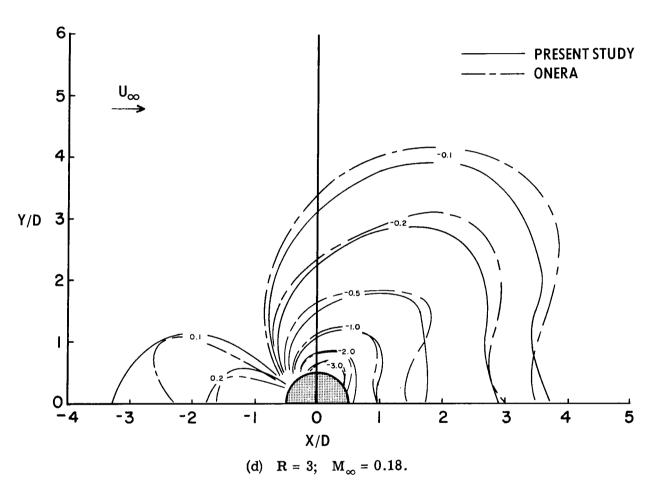
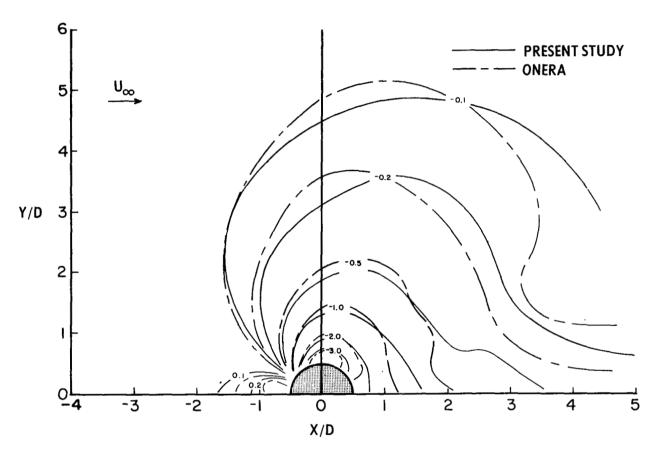
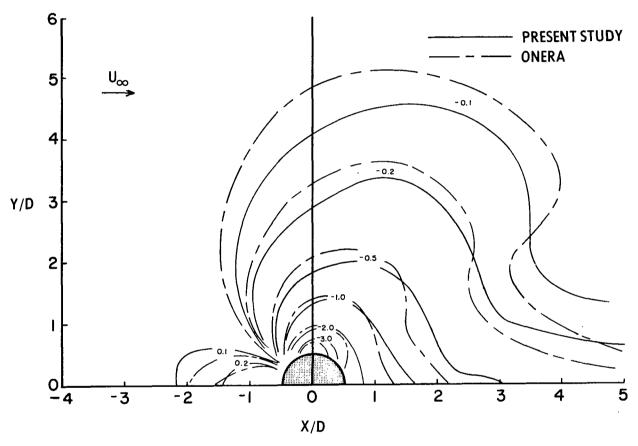


Figure 19.- Continued.



(e) R = 4; $M_{\infty} = 0.09$.

Figure 19.- Continued.



(f) R = 4; $M_{\infty} = 0.18$.

Figure 19.- Continued.

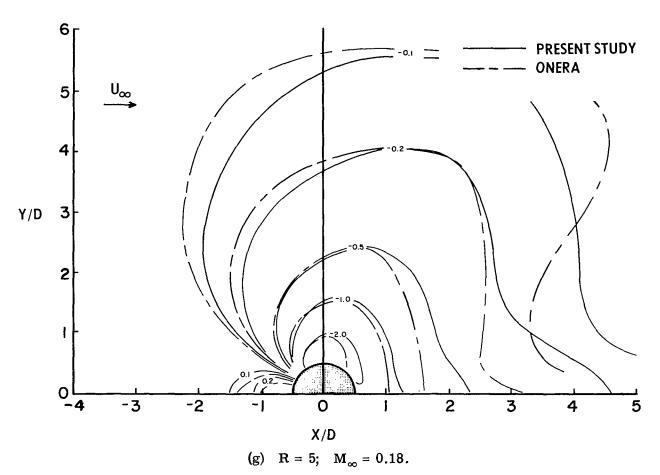
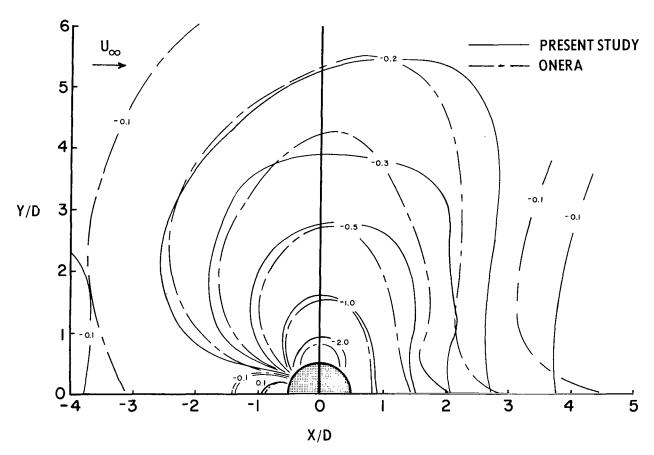


Figure 19.- Continued.



(h) R = 7; $M_{\infty} = 0.09$.

Figure 19.- Continued.

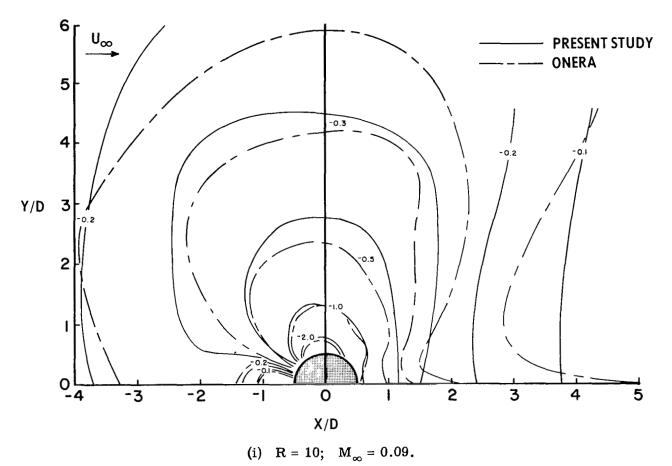


Figure 19.- Concluded.

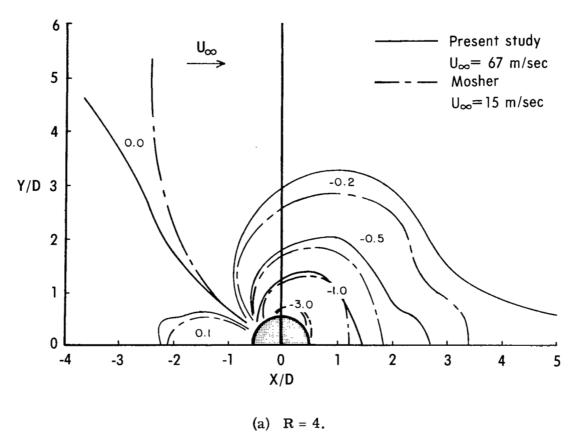


Figure 20.- Comparison with Mosher (ref. 28).

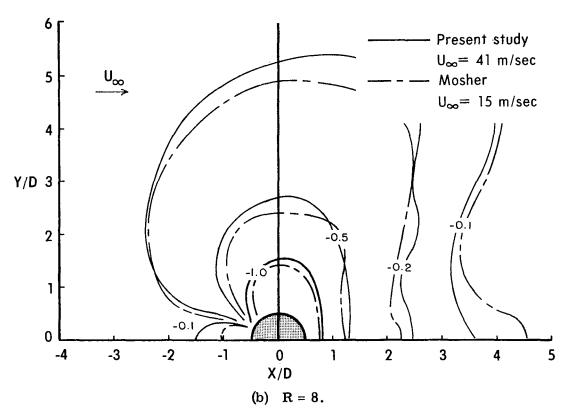


Figure 20.- Concluded.

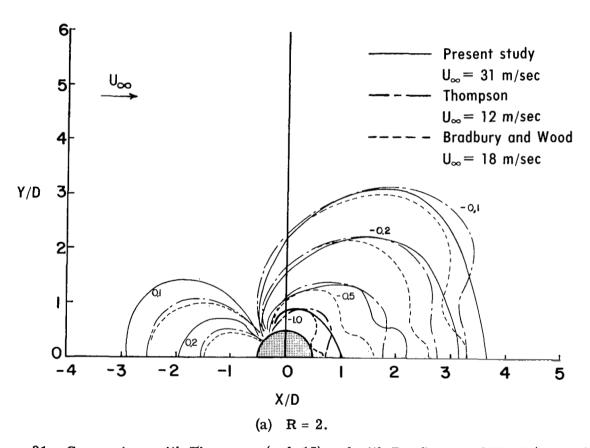


Figure 21.- Comparison with Thompson (ref. 15) and with Bradbury and Wood (ref. 18).

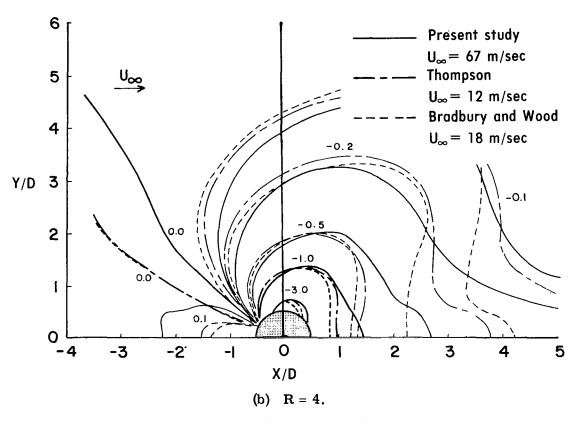


Figure 21.- Continued.

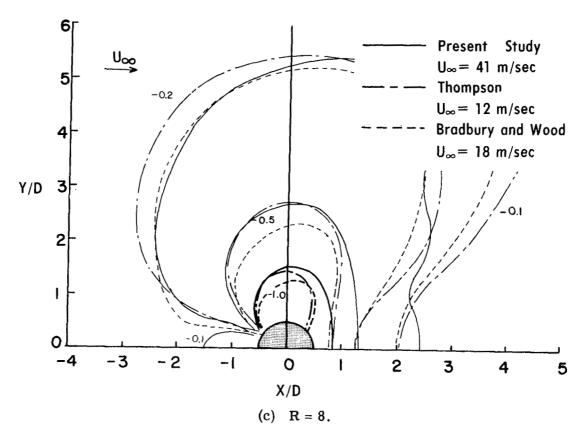


Figure 21.- Concluded.

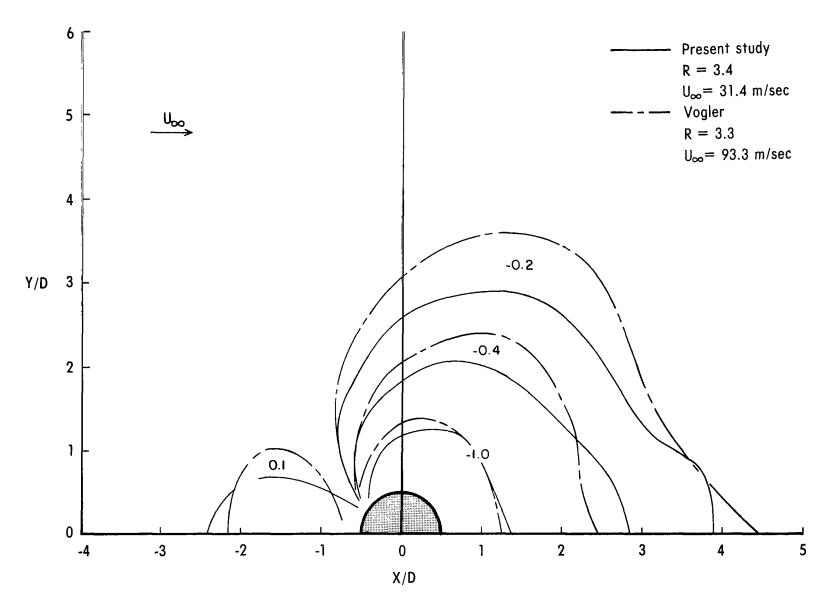


Figure 22.- Comparison with Vogler (ref. 17).

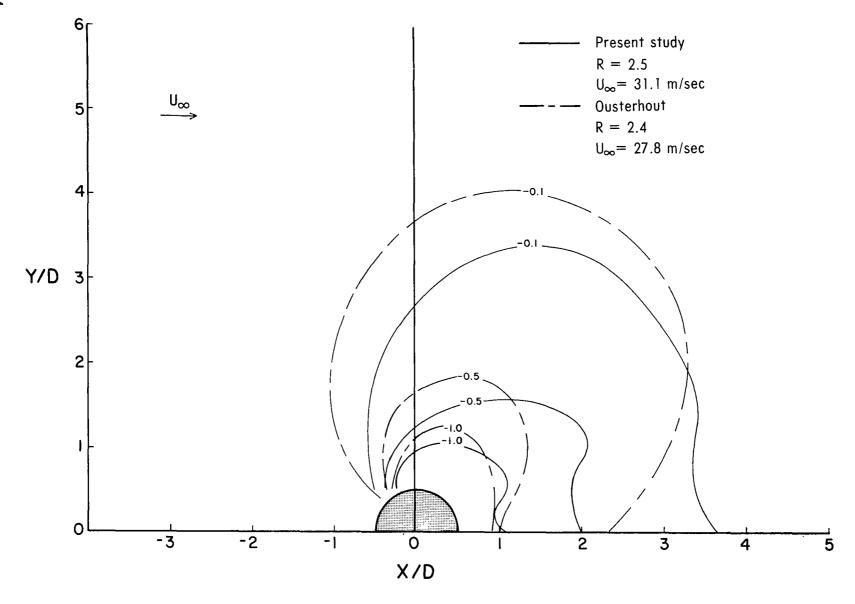


Figure 23.- Comparison with Ousterhout (ref. 23).

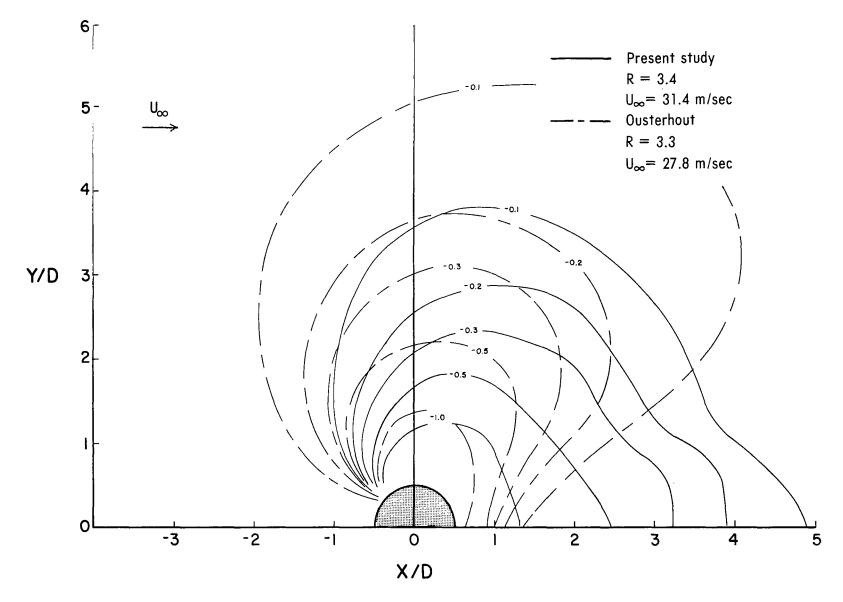


Figure 23.- Concluded.